$$u_t + f(u)_x = 0 \qquad \qquad u = \text{conserved quantity}, \qquad f(u) = \text{flux}$$

$$\frac{d}{dt} \int_{a}^{b} u(t,x) dx = \int_{a}^{b} u_{t}(t,x) dx = -\int_{a}^{b} f(u(t,x))_{x} dx$$
$$= f(u(t,a)) - f(u(t,b)) = [\text{inflow at } a] - [\text{outflow at}]$$



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conservation equation: $u_t + f(u)_x = 0$

quasilinear form: $u_t + f'(u)u_x = 0$

Conservation equation remains meaningful for u = u(t, x) discontinuous, in distributional sense:

$$\iint \left\{ u\phi_t + f(u)\phi_x \right\} \, dxdt = 0 \qquad \text{for all} \quad \phi \in \mathcal{C}^1_c$$

Need only : u, f(u) locally integrable

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$$\mathbf{u_t} + \mathbf{f}(\mathbf{u})_{\mathbf{x}} = \mathbf{0}$$

Assume: u_n is a solution, for every $n \ge 1$,

$$u_n \to u, \qquad f(u_n) \to f(u) \qquad \text{in} \quad \mathbf{L}^1_{loc}$$

then

$$\iint \left\{ u\phi_t + f(u)\phi_x \right\} \, dxdt = \lim_{n \to \infty} \iint \left\{ u_n\phi_t + f(u_n)\phi_x \right\} \, dxdt = 0$$

for all $\phi \in \mathcal{C}^1_c$

(no need to check convergence of derivatives)

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Scalar Conservation Laws

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Scalar Equation with Linear Flux

$$u_t + f(u)_x = 0$$
 $f(u) = \lambda u$

$$u_t + \lambda u_x = 0 \qquad u(0, x) = \phi(x)$$

Explicit solution: $u(t,x) = \phi(x - \lambda t)$



The method of characteristics

$$u_t + f'(u)u_x = 0$$
 $u(0, x) = \phi(x)$

For each x_0 , consider the straight line

$$t \mapsto x(t,x_0) = x_0 + tf'(\phi(x_0))$$

Set $u = \phi(x_0)$ along this line, so that $\dot{x}(t) = f'(u(t, x(t)))$. As long as characteristics do not cross, this yields a solution:



$$u_t + f'(u)u_x = 0$$

Assume: characteristic speed f'(u) is not constant



Global solutions only in a space of discontinuous functions

$$u(t,\cdot)\in BV$$



$$u(t,x) = \begin{cases} u^{-} & \text{if } x < \lambda t \\ u^{+} & \text{if } x > \lambda t \end{cases}$$

is a weak solution if and only if

 $\lambda \cdot [u^+ - u^-] = f(u^+) - f(u^-)$ Rankine - Hugoniot equations

[speed of the shock] \times [jump in the state] = [jump in the flux]

Derivation of the Rankine - Hugoniot equation

$$\iint \left\{ u\phi_t + f(u)\phi_x \right\} dxdt = 0 \quad \text{for all} \quad \phi \in \mathcal{C}_c^1$$

$$\mathbf{v} \doteq \left(u\phi, f(u)\phi \right)$$

$$\mathbf{v} \doteq \left(u\phi, f(u)\phi \right)$$

$$\mathbf{v} \doteq \left(u\phi, f(u)\phi \right)$$

$$\mathbf{v} = \iint_{\Omega^+ \cup \Omega^-} \operatorname{div} \mathbf{v} dxdt = \int_{\partial \Omega^+} \mathbf{n}^+ \cdot \mathbf{v} ds + \int_{\partial \Omega^-} \mathbf{n}^- \cdot \mathbf{v} ds$$

$$= \iint \left[\lambda u^+ - f(u^+) \right] \phi(t, \lambda t) dt + \iint \left[-\lambda u^- + f(u^-) \right] \phi(t, \lambda t) dt$$

$$= \iint \left[\lambda (u^+ - u^-) - \left(f(u^+) - f(u^-) \right) \right] \phi(t, \lambda t) dt$$

$$\lambda (u^{+} - u^{-}) = f(u^{+}) - f(u^{-}) = \int_{0}^{1} f' (\theta u^{+} + (1 - \theta) u^{-}) \cdot (u^{+} - u^{-}) d\theta$$

The Rankine-Hugoniot conditions hold if and only if the speed of the shock is

$$\lambda = \frac{f(u^+) - f(u^-)}{u^+ - u^-} = \int_0^1 f'(\theta u^+ + (1 - \theta)u^-) d\theta$$

= [average characteristic speed]



[speed of the shock] = [slope of secant line through u^- , u^+ on the graph of f]

= [average of the characteristic speeds between u^- and u^+]

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Points of approximate jump

The function u = u(t, x) has an **approximate jump** at a point (τ, ξ) if there exists states $u^- \neq u^+$ and a speed λ such that, calling

$$U(t,x) \doteq \begin{cases} u^- & \text{if } x < \lambda t, \\ u^+ & \text{if } x > \lambda t, \end{cases}$$

there holds



Theorem. If u is a weak solution to a conservation law then the Rankine-Hugoniot equations hold at each point of approximate jump.

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Scalar Conservation Laws



Weak solutions can be non-unique

Example: a Cauchy problem for Burgers' equation

$$u_t + (u^2/2)_x = 0$$
 $u(0,x) = \begin{cases} 1 & \text{if } x \ge 0 \\ 0 & \text{if } x < 0 \end{cases}$

Each $\alpha \in [0, 1]$ yields a weak solution

$$u_{\alpha}(t,x) = \begin{cases} 0 & \text{if} \quad x < \alpha t/2 \\ \alpha & \text{if} \quad \alpha t/2 \le x < (1+\alpha)t/2 \\ 1 & \text{if} \quad x \ge (1+\alpha)t/2 \end{cases}$$



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Stability conditions for shocks

Perturb the shock with left and right states u^- , u^+ by inserting an intermediate state $u^* \in [u^-, u^+]$

Initial shock is stable \iff

 $[{\rm speed \ of \ jump \ behind}] \ \geq \ [{\rm speed \ of \ jump \ ahead}]$



speed of a shock = slope of a secant line to the graph of f



Stability conditions:

- when $u^- < u^+$ the graph of f should remain above the secant line
- when $u^- > u^+$, the graph of f should remain below the secant line

The Lax admissibility condition



A shock connecting the states u^-, u^+ , travelling with speed $\lambda = \frac{f(u^+) - f(u^-)}{u^+ - u^-}$ is admissible if

$$f'(u^-) \geq \lambda \geq f'(u^+)$$

i.e. characteristics do not move out from the shock from either side

Cauchy problem: $u_t + f(u)_x = 0$, $u(0,x) = \bar{u}(x)$ Polygonal approximations of the flux function (Dafermos, 1972) Choose a piecewise affine function f_n such that

$$f_n(u) = f(u)$$
 $u = j \cdot 2^{-n}$, $j \in \mathbb{Z}$

Approximate the initial data with a function $\bar{u}_n : \mathbb{R} \mapsto 2^{-n} \cdot \mathbb{Z}$



Front tracking approximations

piecewise constant approximate solutions: $u_n(t, x)$



 $Tot.Var.(u_n(t, \cdot)) \leq Tot.Var.(\bar{u}_n) \leq Tot.Var.(\bar{u})$

 $\implies \text{ as } n \to \infty, \text{ a subsequence converges in } \mathbf{L}^1_{loc}([0, T] \times \mathbb{R})$ to a weak solution u = u(t, x)

$$u_t + f(u)_x = 0$$

Two initial data in $L^1(\mathbb{R})$: $u_1(0,x) = \overline{u}_1(x), \quad u_2(0,x) = \overline{u}_2(x)$

 ${\bf L}^1$ - distance between solutions does not increase in time:

$$\|u_1(t,\cdot) - u_2(t,\cdot)\|_{\mathsf{L}^1(\mathbb{R})} \leq \|\bar{u}_1 - \bar{u}_2\|_{\mathsf{L}^1(\mathbb{R})}$$

(not true for the L^p distance, p > 1)

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The \mathbf{L}^1 distance between continuous solutions remains constant





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The ${\bf L}^1$ distance decreases when a shock in one solution crosses the graph of the other solution



A related Hamilton-Jacobi equation

$$u_t + f(u)_x = 0$$
 $u(0,x) = \bar{u}(x)$

$$U(t,x) = \int_{-\infty}^{x} u(t,y) \, dy$$

$$U_t + f(U_x) = 0 \qquad \qquad U(0,x) = \overline{U}(x) = \int_{-\infty}^x \overline{u}(y) \, dy$$

f convex ==

U = U(t, x) is the value function for an optimization problem

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$$u \mapsto f(u) \in \mathbb{R} \cup \{+\infty\}$$
 convex

$$f^*(p) \doteq \max_u \{pu - f(u)\}$$



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Scalar Conservation Laws

A representation formula

$$U_t + f(U_x) = 0 \qquad \qquad U(0,x) = \overline{U}(x)$$

$$U(t,x) = \inf_{z(\cdot)} \left\{ \int_0^t f^*(\dot{z}(s)) \, ds + \overline{U}(z(0)); \quad z(t) = x \right\}$$
$$= \min_{y \in \mathbb{R}} \left\{ t \, f^*\left(\frac{x-y}{t}\right) + \overline{U}(y) \right\}$$

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A geometric construction

$$U_{t} + f(U_{x}) = 0 \qquad U(0, x) = \overline{U}(x)$$

define $h(s) \doteq -T f^{*} \left(\frac{-s}{T}\right)$
$$\underbrace{\overline{U}(x)}_{h}$$

$$\underbrace{U(T, x)}_{h} = \inf_{y} \left\{ \overline{U}(y) - h(y - x) \right\}$$

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Cauchy problem:
$$\begin{cases} u_t + f(u)_x = 0, \\ u(0,x) = \overline{u}(x) \end{cases}$$

For each t > 0, and all but at most countably many values of $x \in \mathbb{R}$, there exists a unique y(t, x) s.t.

$$y(t, x) = \arg\min_{y \in \mathbb{R}} \left\{ t f^*\left(\frac{x-y}{t}\right) + \int_{-\infty}^{y} \overline{u}(s) ds \right\}$$

the solution to the Cauchy problem is

$$u(t, x) = (f')^{-1} \left(\frac{x - y(t, x)}{t} \right)$$
(1)



$$y(t, x) = \arg\min_{y \in \mathbb{R}} \left\{ t f^*\left(\frac{x-y}{t}\right) + \int_{-\infty}^{y} \overline{u}(s) ds \right\}$$

define the characteristic speed $\xi \doteq \frac{x - y(t, x)}{t}$

if
$$f'(\omega) = \xi$$
 then $u(t, x) = \omega$

Initial-Boundary value problem



P. Le Floch, Explicit formula for scalar non-linear conservation laws with boundary condition, *Math. Models Appl. Sci.* (1988)

Systems of Conservation Laws

$$\begin{cases} \frac{\partial}{\partial t}u_1 + \frac{\partial}{\partial x}f_1(u_1, \dots, u_n) = 0, \\ & \ddots & \ddots \\ \\ \frac{\partial}{\partial t}u_n + \frac{\partial}{\partial x}f_n(u_1, \dots, u_n) = 0 \end{cases}$$

$$u_t + f(u)_x = 0$$

 $u = (u_1, \ldots, u_n) \in \mathbb{R}^n$ conserved quantities

 $f = (f_1, \ldots, f_n) : \mathbb{R}^n \mapsto \mathbb{R}^n$ fluxes

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Hyperbolic Systems

$$u_t + f(u)_x = 0$$
 $u = u(t, x) \in \mathbb{R}^n$

$$u_t + A(u)u_x = 0 \qquad A(u) = Df(u)$$

The system is **strictly hyperbolic** if each $n \times n$ matrix A(u) has real distinct eigenvalues

$$\lambda_1(u) < \lambda_2(u) < \cdots < \lambda_n(u)$$

right eigenvectors $r_1(u), \ldots, r_n(u)$ (column vectors) left eigenvectors $l_1(u), \ldots, l_n(u)$ (row vectors)

$$Ar_i = \lambda_i r_i \qquad l_i A = \lambda_i l_i$$

Choose bases so that $l_i \cdot r_j = \begin{cases} 1 & \text{if } i = j \\ 0 & \text{if } i \neq j \end{cases}$

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A linear hyperbolic system

$$u_t + Au_x = 0$$
 $u(0, x) = \phi(x)$
 $\lambda_1 < \cdots < \lambda_n$ eigenvalues r_1, \ldots, r_n eigenvectors

Explicit solution: linear superposition of travelling waves



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$$u_t + A(u)u_x = 0$$

eigenvalues depend on $u \implies$ waves change shape



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eigenvectors depend on $u \implies$ nontrivial wave interactions



Global solutions to the Cauchy problem

$$u_t + f(u)_x = 0$$
 $u(0,x) = \bar{u}(x)$

- Construct a sequence of approximate solutions u_m
- Show that (a subsequence) converges: $u_m \rightarrow u$ in L^1_{loc}

 \implies *u* is a weak solution



Need: a-priori bound on the total variation (J. Glimm, 1965)

Building block: the Riemann Problem

$$u_t + f(u)_x = 0$$
 $u(0, x) = \begin{cases} u^- & \text{if } x < 0 \\ u^+ & \text{if } x > 0 \end{cases}$

B. Riemann 1860: 2 × 2 system of isentropic gas dynamics

P. Lax 1957: $n \times n$ systems (+ special assumptions)

T. P. Liu 1975 $n \times n$ systems (generic case)

S. Bianchini 2003 (vanishing viscosity limit for general hyperbolic systems, possibly non-conservative)

invariant w.r.t. symmetry:
$$u^{\theta}(t,x) \doteq u(\theta t, \theta x) \quad \theta > 0$$

Riemann Problem for Linear Systems

$$u_{t} + Au_{x} = 0 \qquad u(0, x) = \begin{cases} u^{-} & \text{if } x < 0\\ u^{+} & \text{if } x > 0 \end{cases}$$

$$u_{t}^{1/1 = \lambda_{2}}$$

$$u^{-} - u^{-} = \sum_{j=1}^{n} c_{j}r_{j} \qquad (\text{sum of eigenvectors of } A)$$

$$u^{+} - u^{-} = \sum_{j=1}^{n} c_{j}r_{j} \qquad (\text{sum of eigenvectors of } A)$$

$$intermediate \text{ states }: \quad \omega_{i} \doteq u^{-} + \sum_{j \leq i} c_{j}r_{j}$$

$$i\text{-th jump: } \omega_{i} - \omega_{i-1} = c_{i}r_{i} \text{ travels with speed } \lambda_{i}$$

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General solution of the Riemann problem: concatenation of elementary waves



Glimm scheme: piecing together solutions of Riemann problems on a fixed grid in the t-x plane



Front tracking scheme: piecing together piecewise constant solutions of Riemann problems at points where fronts interact



$$u_t + f(u)_x = 0, \qquad u(0, x) = \overline{u}(x)$$

Theorem (Glimm 1965).

Assume:

• system is strictly hyperbolic (+ some technical assumptions)

Then there exists $\delta > 0$ such that, for every initial condition $\bar{u} \in L^1(\mathbb{R}; \mathbb{R}^n)$ with

Tot. Var. $(\bar{u}) \leq \delta$,

the Cauchy problem has an entropy admissible weak solution u = u(t, x) defined for all $t \ge 0$.

Uniqueness and continuous dependence on the initial data

$$u_t + f(u)_x = 0$$
 $u(0, x) = \bar{u}(x)$

Theorem (A.B.- R.Colombo, B.Piccoli, T.P.Liu, T.Yang, 1994-1998). For every initial data \bar{u} with small total variation, the front tracking approximations converge to a unique limit solution $u : [0, \infty[\mapsto L^1(\mathbb{R}).$

The flow map $(\bar{u}, t) \mapsto u(t, \cdot) \doteq S_t \bar{u}$ is a uniformly Lipschitz semigroup:

 $S_0 \bar{u} = \bar{u}, \qquad S_s(S_t \bar{u}) = S_{s+t} \bar{u}$

 $\left\|S_t\bar{u}-S_s\bar{v}\right\|_{\mathsf{L}^1} \leq L\cdot\left(\|\bar{u}-\bar{v}\|_{\mathsf{L}^1}+|t-s|\right) \qquad \text{ for all } \bar{u},\bar{v}, \ s,t\geq 0$

Theorem (A.B.- P. LeFloch, M.Lewicka, P.Goatin, 1996-1998).

Any entropy weak solution to the Cauchy problem coincides with the limit of front tracking approximations, hence it is unique

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Scalar Conservation Laws

Claim: weak solutions of the hyperbolic system

$$u_t + f(u)_x = 0$$

can be obtained as limits of solutions to the parabolic system

$$u_t^{\varepsilon} + f(u^{\varepsilon})_x = \varepsilon u_{xx}^{\varepsilon}$$

letting the viscosity $\varepsilon \rightarrow 0+$



Theorem (S. Bianchini, A. Bressan, Annals of Math. 2005)

Consider a strictly hyperbolic system with viscosity

$$u_t + A(u)u_x = \varepsilon \, u_{xx} \qquad \qquad u(0, x) = \bar{u}(x) \,. \tag{CP}$$

If Tot.Var. $\{\bar{u}\}$ is sufficiently small, then (CP) admits a unique solution $u^{\varepsilon}(t, \cdot) = S_t^{\varepsilon} \bar{u}$, defined for all $t \ge 0$. Moreover

$$\mathsf{Tot.Var.}\left\{S_t^{\varepsilon}\bar{u}\right\} \leq C \mathsf{Tot.Var.}\left\{\bar{u}\right\}, \qquad \qquad (\mathsf{BV \ bounds})$$

$$\left\|S_t^{\varepsilon}\bar{u} - S_t^{\varepsilon}\bar{v}\right\|_{\mathsf{L}^1} \leq L \left\|\bar{u} - \bar{v}\right\|_{\mathsf{L}^1} \qquad (\mathsf{L}^1 \text{ stability})$$

(Convergence) If A(u) = Df(u), then as $\varepsilon \to 0$, the viscous solutions u^{ε} converge to the unique entropy weak solution of the system of conservation laws

$$u_t + f(u)_x = 0$$

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- Global existence of solutions to hyperbolic systems for initial data \bar{u} with large total variation
- Existence of entropy weak solutions for systems in **several space dimensions**

