

# Strichartz estimates and sub-Riemannian geometry

## Lecture 2

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## Chapter 3: The Fourier restriction problem

Given a solution  $u(t, x)$  of the classical Schrödinger equation (S) in  $\mathbb{R}^n$

$$\begin{cases} i\partial_t u - \Delta u = 0 \\ u|_{t=0} = u_0, \end{cases}$$

the Fourier transform  $\widehat{u}(t, \xi)$  with respect to the spatial variable  $x$  satisfies

$$i\partial_t \widehat{u}(t, \xi) = -|\xi|^2 \widehat{u}(t, \xi), \quad \widehat{u}(0, \xi) = \widehat{u}_0(\xi). \quad (1)$$

Solving the corresponding ODE and taking the inverse Fourier transform

$$u(t, x) = \int_{\widehat{\mathbb{R}}^n} e^{i(x \cdot \xi + t|\xi|^2)} \widehat{u}_0(\xi) d\xi. \quad (2)$$

One can also interpret it as the inverse Fourier transform of a data on the paraboloid  $\widehat{S}$  in the space of frequencies

$$u(t, x) = \int_{\mathbb{R}^n} e^{i(x \cdot \xi + t|\xi|^2)} \widehat{u}_0(\xi) d\xi = \int_{\widehat{S}} e^{iy \cdot z} g(z) d\sigma(z)$$

where  $\widehat{\mathbb{R}}^{n+1} = \widehat{\mathbb{R}} \times \widehat{\mathbb{R}}^n$ , defined as

$$\widehat{S} \stackrel{\text{def}}{=} \{(\alpha, \xi) \in \widehat{\mathbb{R}} \times \widehat{\mathbb{R}}^n \mid \alpha = |\xi|^2\}.$$

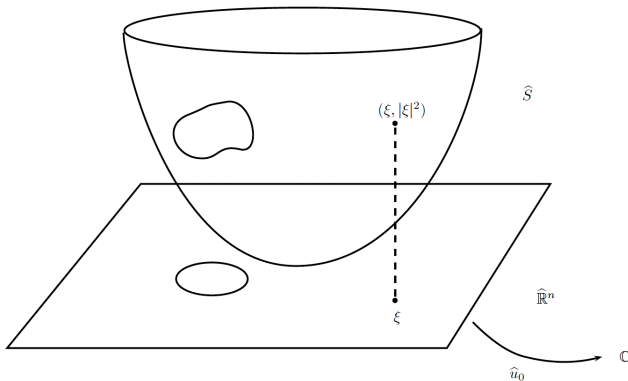
where  $y = (t, x)$  and  $z = (\alpha, \xi)$

$$\|u\|_{L^{p'}(\mathbb{R}^{n+1})} = \|\mathcal{F}^{-1}(gd\sigma)\|_{L^{p'}(\mathbb{R}^{n+1})}$$

# Geometric interpretation



- Let us endow  $\widehat{S}$  with the measure  $d\sigma = d\xi$ .
- $d\sigma$  is not the intrinsic surface measure of  $\widehat{S}$ , which is  $d\mu = \sqrt{1 + 2|\xi|}d\xi$ .



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## RESTRICTIONS OF FOURIER TRANSFORMS TO QUADRATIC SURFACES AND DECAY OF SOLUTIONS OF WAVE EQUATIONS

ROBERT S. STRICHARTZ

### §1. Introduction

Let  $S$  be a subset of  $\mathbb{R}^n$  and  $d\mu$  a positive measure supported on  $S$  and of temperate growth at infinity. We consider the following two problems:

*Problem A.* For which values of  $p$ ,  $1 \leq p < 2$ , is it true that  $f \in L^p(\mathbb{R}^n)$  implies  $\hat{f}$  has a well-defined restriction to  $S$  in  $L^2(d\mu)$  with

$$(1.1) \quad \left( \int |\hat{f}|^2 d\mu \right)^{1/2} \leq c_p \|f\|_p?$$

*Problem B.* For which values of  $q$ ,  $2 < q \leq \infty$ , is it true that the tempered distribution  $Fd\mu$  for each  $F \in L^2(d\mu)$  has Fourier transform in  $L^q(\mathbb{R}^n)$  with

$$(1.2) \quad \|(Fd\mu)^\wedge\|_q \leq c_q \left( \int |F|^2 d\mu \right)^{1/2}?$$

A simple duality argument shows these two problems are completely equivalent if  $p$  and  $q$  are dual indices,  $(1/p) + (1/q) = 1$ . Interest in Problem A when  $S$  is a sphere stems from the work of C. Fefferman [3], and in this case the answer is known (see [11]). Interest in Problem B was recently signalled by I. Segal [6] who studied the special case  $S = \{(x, y) \in \mathbb{R}^2 : y^2 - x^2 = 1\}$  and gave the interpretation of the answer as a space-time decay for solutions of the Klein-Gordon equation with finite relativistic-invariant norm.

In this paper we give a complete solution when  $S$  is a quadratic surface given by

$$(1.3) \quad S = \{x \in \mathbb{R}^n : R(x) = r\}$$

where  $R(x)$  is a polynomial of degree two with real coefficients and  $r$  is a real constant. To avoid triviality we assume  $R$  is not a function of fewer than  $n$  variables, so that aside from isolated points  $S$  is a  $n - 1$ -dimensional  $C^\infty$  manifold. There is a canonical measure  $d\mu$  associated to the function  $R$  (not intrinsic to the surface  $S$ , however) given by

$$(1.4) \quad d\mu = \frac{dx_1 \cdots dx_{n-1}}{|\partial R / \partial x_n|}$$

A lot of contributors: Stein, Fefferman, Tomas, etc.

**Problem:** Can we restrict Fourier transform of  $L^p$  functions to subsets ?

- $f$  in  $L^1(\mathbb{R}^n)$  implies  $\mathcal{F}(f)$  continuous  $\rightarrow$  OK.
- $f$  in  $L^2(\mathbb{R}^n)$  implies  $\mathcal{F}(f)$  in  $L^2(\widehat{\mathbb{R}}^n)$   $\rightarrow$  arbitrary on a zero meas set  $\widehat{S}$  of  $\widehat{\mathbb{R}}^n$ .
- what happens for  $1 < p < 2$ ?
- it depends on the surface!
- if the surface is “flat” we cannot do a lot



- The Fourier transform of a  $L^p$  function, for **any**  $p > 1$ , **cannot** be restricted to hyperplanes.
- This  $f$  belongs to  $L^p(\mathbb{R}^n)$ , for all  $p > 1$

$$f(x) = \frac{e^{-|x'|^2}}{1 + |x_1|} \quad x = (x_1, x') \in \mathbb{R}^n, \quad (3)$$

- its Fourier transform does not admit a restriction on  $\widehat{S} = \{\xi_1 = 0\}$ .

$$\widehat{f}(0, \xi') = \int_{\mathbb{R}^n} e^{-ix' \cdot \xi'} \frac{e^{-|x'|^2}}{1 + |x_1|} dx_1 dx'$$

→ what happens for different surfaces?

## Tomas and Stein

One can restrict the Fourier transform of  $L^p(\mathbb{R}^n)$  functions, for  $p > 1$  (close to 1), to hypersurfaces  $\widehat{S}$  that are “sufficiently curved”, (main example: the sphere).

Let us state more formally the questions

**Problem:** given a hypersurface  $\widehat{S} \subset \widehat{\mathbb{R}}^n$  endowed with a smooth measure  $d\sigma$ , the restriction problem asks for which pairs  $(p, q)$  an inequality of the form

$$\|\mathcal{F}(f)|_{\widehat{S}}\|_{L^q(\widehat{S}, d\sigma)} \leq C \|f\|_{L^p(\mathbb{R}^n)} \quad (4)$$

holds for all  $f$  in  $\mathcal{S}(\mathbb{R}^n)$ .

- The operator  $R_S$  is continuous from  $L^p(\mathbb{R}^n)$  to  $L^q(\widehat{S}, d\sigma)$ ?

$$R_S f = \mathcal{F}(f)|_{\widehat{S}}$$

→ not completely settled in its general form

from now on

we focus on the case  $q = 2$

- the adjoint operator  $R_S^*$  is continuous from  $L^2(\widehat{S}, d\sigma)$  to  $L^{p'}(\mathbb{R}^n)$ ?

$$R_S^* g = \mathcal{F}^{-1}(g d\sigma)$$

$$\|\mathcal{F}^{-1}(g d\sigma)\|_{L^{p'}(\mathbb{R}^n)} \leq C \|g\|_{L^2(\widehat{S}, d\sigma)} \quad (5)$$

A basic counterexample shows that the range of  $p$  for which the estimate holds cannot be the entire interval  $1 \leq p \leq 2$ ;

## Example (Knapp)

Let  $\widehat{S}$  be the  $(n-1)$ -dimensional sphere in  $\widehat{\mathbb{R}}^n$  endowed with the standard measure  $d\mu$ . The estimate can hold only if  $p \leq \frac{2n+2}{n+3} = 2 - \frac{4}{n+3}$ .

- Consider the equivalent formulation of the estimate

$$\|\widehat{g\sigma}\|_{L^{p'}(\mathbb{R}^n)} \leq C \|g\|_{L^2(S^{n-1})} \quad (6)$$

- Let  $\delta > 0$  and let  $g_\delta$  be the characteristic function “spherical cap”

$$\widehat{C}_\delta = \{x \in \widehat{S} : |x \cdot e_n| < \delta\}.$$

- We consider the equivalent formulation of estimate

$$\|\widehat{g\sigma}\|_{L^{p'}(\mathbb{R}^n)} \leq C \|g\|_{L^2(S^{n-1})} \quad (7)$$

- Let  $\delta > 0$  be small and let  $g_\delta$  be the characteristic function on  $C_\delta$ .
- $|C_\delta| \sim \delta^{n-1}$ . This implies  $\|g_\delta\|_{L^2(S^{n-1})} \sim \delta^{(n-1)/2}$ .
- If  $x \in \mathbb{R}^n$  is orthogonal to the vertical direction

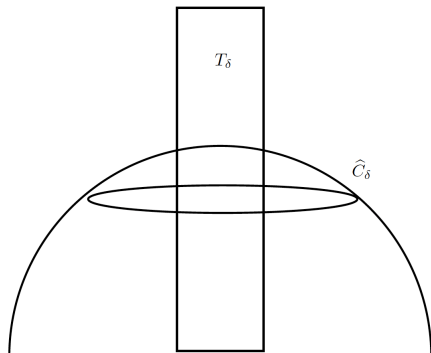
$$|\widehat{g_\delta\sigma}(x)| = \left| \int_{S^{n-1}} e^{ix \cdot \xi} g_\delta(\xi) d\sigma(\xi) \right| = \left| \int_{C_\delta} e^{ix \cdot \xi} d\sigma(\xi) \right| \sim |C_\delta| \sim \delta^{n-1}.$$

$$\|\widehat{g_\delta\sigma}\|_{L^{p'}(\mathbb{R}^n)} = \left( \int_{\mathbb{R}^n} |\widehat{g_\delta\sigma}(x)|^{p'} dx \right)^{1/p'}$$

- let  $T_\delta$  be the tube in the  $x$  space oriented orthogonally to the sphere

$$[-\delta^{-1}, \delta^{-1}] \times \dots \times [-\delta^{-1}, \delta^{-1}] \times [-\delta^{-2}, \delta^{-2}]$$

- $|T_\delta| \sim \delta^{-n-1}$ .



- For  $x$  in  $T_\delta$  and  $\delta$  very small the quantity  $x \cdot \xi$  is almost zero for  $\xi \in C_\delta$ .

$$\|\widehat{g_\delta \sigma}\|_{L^{p'}(\mathbb{R}^n)} = \left( \int_{\mathbb{R}^n} |\widehat{g_\delta \sigma}(x)|^{p'} dx \right)^{1/p'} \quad (8)$$

$$\geq \left( \int_{T_\delta} |\widehat{g_\delta \sigma}(x)|^{p'} dx \right)^{1/p'} \quad (9)$$

$$\sim \left( \int_{T_\delta} \delta^{(n-1)p'} dx \right)^{1/p'} \quad (10)$$

$$\sim \delta^{(n-1)} |T_\delta|^{1/p'} \sim \delta^{(n-1)} \delta^{(-n-1)/p'} \quad (11)$$

The estimate can hence be valid only if (the inequality is  $\geq$  since  $\delta \rightarrow 0$ )

$$n - 1 - \frac{n + 1}{p'} \geq \frac{n - 1}{2}$$

which is the conclusion.

The above range is indeed the correct one for non vanishing curvature.

## Theorem (Tomas-Stein, 1975)

Let  $\widehat{S}$  be a smooth compact hypersurface in  $\widehat{\mathbb{R}}^n$  with non vanishing Gaussian curvature at every point, and let  $d\sigma$  be a smooth measure on  $\widehat{S}$ . Then

$$\|\mathcal{F}(f)|_{\widehat{S}}\|_{L^2(\widehat{S}, d\sigma)} \leq C_p \|f\|_{L^p(\mathbb{R}^n)}.$$

for every  $f \in \mathcal{S}(\mathbb{R}^n)$  and every  $p \leq (2n + 2)/(n + 3)$ ,

- A similar result is possible for surfaces with vanishing Gaussian curvature (that are not flat).
- In this case the range of  $p$  is smaller depending on the order of tangency of the surface to its tangent space.
- The assumption about compactness of  $\widehat{S}$  can be removed by replacing  $d\sigma$  with a compactly supported smooth measure.



Equivalent to the continuity from  $L^p(\mathbb{R}^n)$  to  $L^{p'}(\mathbb{R}^n)$  of the operator

$$R_S^* R_S f = f * \hat{\sigma} \quad (12)$$

$$\|\mathcal{F}(f)|_{\hat{S}}\|_{L^2(\hat{S}, d\sigma)}^2 = \int (f * \hat{\sigma}) f dx \leq \|f * \hat{\sigma}\|_{L^{p'}(\mathbb{R}^n)} \|f\|_{L^p(\mathbb{R}^n)}$$

Recall that the Fourier transform of the measure  $d\sigma$  is a function given by

$$\hat{\sigma}(\xi) = \int_{\mathbb{R}^n} e^{ix \cdot \xi} d\sigma(x) \quad (13)$$

Let  $S$  be a smooth compact hypersurfaces with non-zero Gaussian curvature at every point. Then

$$|\hat{\sigma}(\xi)| \leq C(1 + |\xi|)^{-\frac{n-1}{2}} \quad (14)$$

Let  $S$  be a smooth compact hypersurfaces with non-zero Gaussian curvature at every point. Then

$$|\widehat{\sigma}(\xi)| \leq C(1 + |\xi|)^{-\frac{n-1}{2}} \quad (15)$$

- only with decay one only gets  $p \leq \frac{4n}{3n+1}$  (Fefferman, Stein)

$$n = 3, \quad \widehat{\sigma}(\xi) = 2 \frac{\sin(2\pi|x|)}{|x|}$$

- using a dyadic decomposition and real interpolation  $p < \frac{2(n+1)}{n+3}$  (Tomas)
- with complex interpolation  $p = \frac{2(n+1)}{n+3}$  (Stein)

The classical Schrödinger equation in  $\mathbb{R}^n$ : taking the inverse Fourier transform

$$u(t, x) = \int_{\widehat{\mathbb{R}}^n} e^{i(x \cdot \xi + t|\xi|^2)} \widehat{u}_0(\xi) d\xi. \quad (16)$$

Consider the paraboloid  $\widehat{S}$  in the space of frequencies  $\widehat{\mathbb{R}}^{n+1} = \widehat{\mathbb{R}} \times \widehat{\mathbb{R}}^n$

$$\widehat{S} = \left\{ (\alpha, \xi) \in \widehat{\mathbb{R}} \times \widehat{\mathbb{R}}^n \mid \alpha = |\xi|^2 \right\}.$$

- Given  $\widehat{u}_0 : \widehat{\mathbb{R}}^n \rightarrow \mathbb{C}$  define  $g : \widehat{S} \rightarrow \mathbb{C}$  as  $g(|\xi|^2, \xi) = \widehat{u}_0(\xi)$ . Then

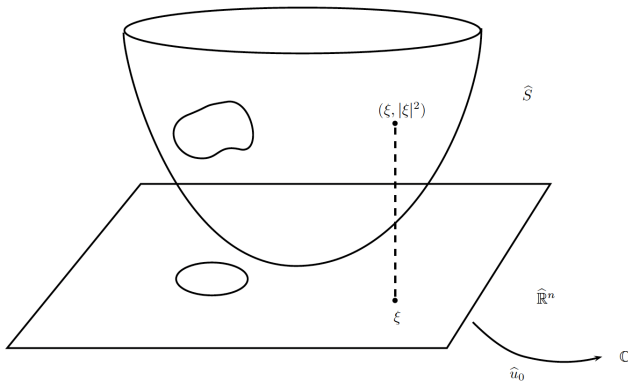
$$u(t, x) = \int_{\widehat{\mathbb{R}}^n} e^{i(x \cdot \xi + t|\xi|^2)} \widehat{u}_0(\xi) d\xi = \int_{\widehat{S}} e^{iy \cdot z} g(z) d\sigma(z)$$

where  $y = (t, x)$  and  $z = (\alpha, \xi)$ .

# Geometric interpretation



- Let us endow  $\widehat{S}$  with the measure  $d\sigma = d\xi$ .
- $d\sigma$  is not the intrinsic surface measure of  $\widehat{S}$ , which is  $d\mu = \sqrt{1 + 2|\xi|^2}d\xi$ .



## The Fourier restriction theorem

$$\|\mathcal{F}^{-1}(gd\sigma)\|_{L^{p'}(\widehat{\mathbb{R}^{n+1}})} \leq C_p \|g\|_{L^2(\widehat{S}, d\mu)}, \quad (17)$$

for all  $g \in L^2(\widehat{S}, d\mu)$  and all  $p' \geq 2(n+2)/n$ .

By construction  $\|g\|_{L^2(\widehat{S}, d\mu)} = \|\widehat{u}_0\|_{L^2(\widehat{\mathbb{R}^n})} = \|u_0\|_{L^2(\mathbb{R}^n)}$

→ we stress that we apply the result in dimension  $n+1$ , i.e., in  $\mathbb{R} \times \mathbb{R}^n = \mathbb{R}^{n+1}$

Applying the statement to  $g$  related to a initial data  $u_0$  such that  $\widehat{u}_0$  is supported on a unit ball

$$\|u\|_{L^{p'}(\mathbb{R}^{n+1})} \leq C \|u_0\|_{L^2(\mathbb{R}^n)}, \quad (18)$$

for all  $p' \geq 2(n+2)/n$ .

A scaling argument and the density of spectrally localized functions in  $L^2(\mathbb{R}^n)$ , give the result for  $p' = 2 + \frac{4}{n}$ , and all  $p' \geq 2(n+2)/n$ .

1. Prove a Fourier restriction on the Heisenberg group
  - a result of D.Müller  $\rightarrow$  specific for the sphere
  - what is the sphere? what about paraboloid?
2. We do not exactly need restriction theorems for  $\mathbb{H}^d$ 
  - we applied the result to a surface in the space  $\mathbb{R}^{n+1} = \mathbb{R} \times \mathbb{R}^n$ $\rightarrow$  the paraboloid for the Schrödinger eq. (the cone for the wave equation).
  - when dealing with equations defined on the Heisenberg group  $\mathbb{H}^d$ , one is naturally lead to consider surfaces in the space  $\mathbb{R} \times \widehat{\mathbb{H}}^d$ , which is not related to  $\mathbb{H}^{d'}$  for some  $d'$ .

## Chapter 4: Strichartz estimates in the Heisenberg group

A function  $\phi$  on  $\mathbb{H}^1$  is said to be *radial* if  $\phi(x, y, z) = \phi(x^2 + y^2, z)$ .

## Theorem (Bahouri, DB, Gallagher, '21)

Given  $(p, q)$  belonging to the admissible set

$$\mathcal{A} = \left\{ (p, q) \in [2, \infty]^2 / p \leq q \quad \text{and} \quad \frac{2}{q} + \frac{2d}{p} = \frac{Q}{2} \right\},$$

the solution to the Schrödinger equation ( $S_{\mathbb{H}}$ ) with radial data satisfies

$$\|u\|_{L_z^\infty L_t^q L_{x,y}^p} \leq C_{p,q,p_1,q_1} \left( \|u_0\|_{L^2(\mathbb{H}^d)} \right).$$

- restrictive due to  $p \leq q$ . Indeed  $p = q = 2$ .
- we stress that  $L_z^\infty L_t^q L_{x,y}^p \neq L_t^\infty L_z^q L_{x,y}^p$
- similar for inhomogeneous and wave



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the solution to the Schrödinger equation  $(S_{\mathbb{H}})$  with radial data satisfies

$$\|u\|_{L_z^\infty L_t^q L_{x,y}^p} \leq C_{p,q,p_1,q_1} \left( \|u_0\|_{H^\sigma(\mathbb{H}^d)} \right).$$

- $\sigma = \frac{Q}{2} - \frac{2}{q} - \frac{2d}{p}$  is the loss of derivatives,  $\sigma = 0$  forces  $p = q$
- we stress that  $L_z^\infty L_t^q L_{x,y}^p \neq L_t^\infty L_z^q L_{x,y}^p$
- similar for inhomogeneous and wave

A function  $\phi$  on  $\mathbb{H}^d$  is said to be *radial* if  $\phi(z, s) = f(|z|, s)$ .

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Given  $(p, q)$  and  $(p_1, q_1)$  belonging to the admissible set

$$\mathcal{A} = \left\{ (p, q) \in [2, \infty]^2 / q \leq p \quad \text{and} \quad \frac{2}{q} + \frac{2d}{p} \leq \frac{Q}{2} \right\},$$

the solution to the Schrödinger equation ( $S_{\mathbb{H}}$ ) with radial data satisfies

$$\|u\|_{L_s^\infty L_t^q L_z^p} \leq C_{p,q,p_1,q_1} \left( \|u_0\|_{H^\sigma(\mathbb{H}^d)} + \|f\|_{L_t^1 H^\sigma(\mathbb{H}^d)} \right).$$

It is defined using **irreducible unitary representations** : for any integrable function  $u$  on  $\mathbb{H}$  (Kirillov theory)

$$\forall \lambda \in \mathbb{R}^*, \quad \widehat{u}(\lambda) := \int_{\mathbb{H}} u(x) \mathcal{R}_x^\lambda dx,$$

with  $\mathcal{R}^\lambda$  the group homomorphism between  $\mathbb{H}$  and the unitary group  $\mathcal{U}(L^2(\mathbb{R}))$  of  $L^2(\mathbb{R})$  given for all  $x$  in  $\mathbb{H}$  and  $\phi$  in  $L^2(\mathbb{R})$ , by

$$\mathcal{R}_x^\lambda \phi(\theta) := \exp\left(i\lambda x_3 + i\lambda \theta x_2\right) \phi(\theta + x_1).$$

Then  $\widehat{u}(\lambda)$  is a family of bounded operators on  $L^2(\mathbb{R})$ , with many properties similar to  $\mathbb{R}^d$  : inversion formula, Fourier-Plancherel identity  
*Trace* *Hilbert – Schmidt*

The sub-Laplacian

$$\Delta_{\mathbb{H}} = X_1^2 + X_2^2$$

There holds

$$\widehat{-\Delta_{\mathbb{H}}u}(\lambda) = \widehat{u}(\lambda) \circ P_{\lambda}, \quad \text{with} \quad P_{\lambda} := -\frac{d^2}{d\theta^2} + \lambda^2\theta^2.$$

The spectrum of the **rescaled harmonic oscillator** is

$$\text{Sp}(P_{\lambda}) = \{|\lambda|(2m+1), m \in \mathbb{N}\}$$

and the eigenfunctions are the Hermite functions  $\psi_m^{\lambda}$ . So for all  $m \in \mathbb{N}$ ,

$$\widehat{-\Delta_{\mathbb{H}}u}(\lambda)\psi_m^{\lambda} = E_m(\lambda)\widehat{u}(\lambda)\psi_m^{\lambda}.$$

Set  $\widehat{x} := (n, m, \lambda) \in \widehat{\mathbb{H}} = \mathbb{N}^2 \times \mathbb{R}^*$ , and

$$\begin{aligned}\mathcal{F}_{\mathbb{H}}(u)(n, m, \lambda) &:= (\widehat{u}(\lambda)\psi_m^\lambda|\psi_n^\lambda)_{L^2(\mathbb{R})} \\ &= \int_{\mathbb{H}} \mathcal{W}(\widehat{x}, x)u(x)dx\end{aligned}$$

where  $\mathcal{W}(\widehat{x}, x) := e^{i\lambda x_3} e^{-|\lambda|(x_1^2+x_2^2)} \underbrace{L_m(2|\lambda|(x_1^2+x_2^2))}_{\text{Laguerre polynomial}}$ .

Then

$$\mathcal{F}_{\mathbb{H}}(-\Delta_{\mathbb{H}}u)(n, m, \lambda) = \underbrace{E_m(\lambda)}_{\text{frequency}} \mathcal{F}_{\mathbb{H}}(u)(n, m, \lambda).$$

Bahouri, Chemin, Danchin

Inversion and Fourier-Plancherel formulae

$$f(\widehat{x}) = \frac{2^{d-1}}{\pi^{d+1}} \int_{\widetilde{\mathbb{H}}^d} \mathcal{W}(\widehat{x}, x) \mathcal{F}_{\mathbb{H}} f(\widehat{x}) d\widehat{x}$$

and

$$(\mathcal{F}_{\mathbb{H}} f | \mathcal{F}_{\mathbb{H}} g)_{L^2(\widetilde{\mathbb{H}}^d)} = \frac{\pi^{d+1}}{2^{d-1}} (f | g)_{L^2(\mathbb{H}^d)},$$

Action of the Laplacian

$$\mathcal{F}_{\mathbb{H}}(\Delta_{\mathbb{H}} f)(\widehat{x}) = -4|\lambda|(2|m| + d) \mathcal{F}_{\mathbb{H}}(f)(\widehat{x}).$$

Radial functions  $f(z, s) = f(|z|, s)$

$$\mathcal{F}_{\mathbb{H}}(f)(n, m, \lambda) = \mathcal{F}_{\mathbb{H}}(f)(n, m, \lambda) \delta_{n,m} = \mathcal{F}_{\mathbb{H}}(f)(|n|, |n|, \lambda) \delta_{n,m}.$$

Convolution for radial functions

$$\mathcal{F}_{\mathbb{H}}(f \star g)(\ell, \ell, \lambda) = \mathcal{F}_{\mathbb{H}} f(\ell, \ell, \lambda) \mathcal{F}_{\mathbb{H}} g(\ell, \ell, \lambda)$$

Let  $u_0$  in  $\mathcal{S}(\mathbb{H}^d)$  be **radial** and consider the Cauchy problem

$$\begin{cases} i\partial_t u - \Delta_{\mathbb{H}} u = 0 \\ u|_{t=0} = u_0. \end{cases}$$

Taking the partial Fourier transform with respect to the variable  $w$

$$\begin{cases} i\frac{d}{dt}\mathcal{F}_{\mathbb{H}}(u)(t, n, m, \lambda) = -4|\lambda|(2|m| + d)\mathcal{F}_{\mathbb{H}}(u)(t, n, m, \lambda) \\ \mathcal{F}_{\mathbb{H}}(u)|_{t=0} = \mathcal{F}_{\mathbb{H}}u_0. \end{cases}$$

$$\mathcal{F}_{\mathbb{H}}(u)(t, n, m, \lambda) = e^{4it|\lambda|(2|m|+d)}\mathcal{F}_{\mathbb{H}}(u_0)(|n|, |n|, \lambda)\delta_{n,m}.$$

→ Notice that if we set  $|m| = 0$  we see the “transport” part

$$\mathcal{F}_{\mathbb{H}}(u)(t, 0, 0, \lambda) = e^{4it|\lambda|d}\mathcal{F}_{\mathbb{H}}(u_0)(0, 0, \lambda).$$

Applying the inverse Fourier formula

$$u(t, z, s) = \frac{2^{d-1}}{\pi^{d+1}} \int_{\widehat{\mathbb{H}}^d} \mathcal{W}(\widehat{x}, z, s) e^{4it|\lambda|(2|m|+d)} \mathcal{F}_{\mathbb{H}}(u_0)(|n|, |n|, \lambda) \delta_{n,m} d\widehat{x}.$$

Re-expressed as the inverse Fourier transform in  $\widehat{\mathbb{R}} \times \widehat{\mathbb{H}}^d$  of  $\mathcal{F}_{\mathbb{H}}(u_0) d\Sigma$ ,

$$\Sigma \stackrel{\text{def}}{=} \left\{ (\alpha, \widehat{x}) = (\alpha, (n, n, \lambda)) \in \widehat{\mathbb{R}} \times \widehat{\mathbb{H}}^d / \alpha = 4|\lambda|(2|n| + d) \right\}.$$

endow  $\Sigma$  with the measure  $d\Sigma$  induced by the projection  $\widehat{\mathbb{R}} \times \widehat{\mathbb{H}}^d \rightarrow \widehat{\mathbb{H}}^d$

$$\int_{\widehat{\mathbb{D}}} \Phi(\alpha, \widehat{x}) d\Sigma(\alpha, \widehat{x}) = \int_{\widehat{\mathbb{H}}^d} \Phi(4|\lambda|(2|m| + d), \widehat{x}) d\widehat{x},$$

Theorem (Bahouri, DB, Gallagher, '19)

If  $1 \leq q \leq p \leq 2$ , then for  $f$  radial

$$\|\mathcal{F}_{\widehat{\mathbb{R}} \times \widehat{\mathbb{H}}^d}(f)|_{\Sigma}\|_{L^2(d\Sigma)} \leq C_{p,q} \|f\|_{L_s^1 L_t^q L_z^p}, \quad (19)$$



Using dual inequality, assuming that  $\mathcal{F}_{\mathbb{H}} u_0$  is localized in the unit ball

For any  $2 \leq p \leq q \leq \infty$

$$\|u\|_{L_s^\infty L_t^q L_z^p} \leq C \|\mathcal{F}_{\mathbb{H}} u_0\|_{L^2(\widehat{\mathbb{H}}^d)} = C \|u_0\|_{L^2(\mathbb{H}^d)},$$

- If  $u_0$  is frequency localized in the ball  $\mathcal{B}_\Lambda$ ,

$$u_\Lambda(t, z, s) = u(\Lambda^{-2}t, \Lambda^{-1}z, \Lambda^{-2}s), \quad u_{0,\Lambda}(z, s) = u_0(\Lambda^{-1}z, \Lambda^{-2}s)$$

- we have

$$\|u_\Lambda\|_{L_s^\infty L_t^q L_z^p} = \Lambda^{\frac{2}{q} + \frac{2d}{p}} \|u\|_{L_s^\infty L_t^q L_z^p}, \quad \|u_{0,\Lambda}\|_{L^2(\mathbb{H}^d)} = \Lambda^{\frac{Q}{2}} \|u_0\|_{L^2(\mathbb{H}^d)},$$

- we infer

$$\|u\|_{L_s^\infty L_t^q L_z^p} \leq C \Lambda^{\frac{Q}{2} - \frac{2}{q} - \frac{2d}{p}} \|u_0\|_{L^2(\mathbb{H}^d)}.$$