

THE CONTINUITY EQUATION

Consider $E(t) \subset \bar{E}$, $t \in I$, τ measurable
 $E(t)$ measurable $\forall t$

$\partial E(t)$ be of class C^1 (uniformly in t)

Consider $\rho \in C^1(E \times I)$

We want to compute $\frac{d}{dt} \int_{E(t)} \rho(x,t) dx$

Given a family of $n \times n$ matrices $t \mapsto A(t)$

depending on a parameter $t \in I$,

$t \mapsto A(t) \in C^1(I)$ (i.e. $(A(t))_{ij} = a_{ij}(t)$)

and $t \mapsto a_{ij}(t) \in C^1(I)$

$\forall i, j$

one has that

$$\frac{d}{dt} (\det A(t)) = \det A(t) \cdot \text{tr} \left(A^{-1}(t) \frac{d}{dt} A(t) \right)$$

$$\text{where } \left(\frac{d}{dt} A(t) \right)_{ij} = \frac{d}{dt} a_{ij}(t)$$

proof: the result is given without proof, even if for the sake of completeness we put it at the end of the section.

We now go on computing $\frac{d}{dt} \int_{E(t)} p(x,t) dx$.

Assumptions

We assume that $E(t) = G_t(E_0)$

$$G_t(x) = G(t,x) \quad G_t \text{ bijection uniformly } C^1$$

$$G_t^{-1} \text{ uniformly } C^1$$

$G \in C^1$ with respect to t and

$$|DG(x, \cdot)| = |\det JG_t(x)| \quad C^1 \text{ in } t$$

$$E(t) = G_t(E_0)$$

$$\frac{d}{dt} \int_{E(t)} p(y,t) dy = \left(y = G_t(x) \right)$$

$$= \frac{d}{dt} \int_{E_0} p(G_t(x), t) |DG_t(x)| dx =$$

E_0

$$= \int_{E_0} \frac{\partial}{\partial t} \left(p(G_t(x), t) |DG_t(x)| \right) dx =$$

$$\begin{aligned}
&= \int_{E_0} \sum_{i=1}^n \frac{\partial \ell}{\partial y_i} (G_t(x), t) \frac{\partial (G_t)_i}{\partial t} (x) |DG_t(x)| dx + \\
&\quad + \int_{E_0} \frac{\partial \ell}{\partial t} (G_t(x), t) |DG_t(x)| dx + \\
&\quad + \int_{E_0} \rho (G(x,t), t) \cdot \frac{\partial}{\partial t} |DG(x,t)| dx
\end{aligned}$$

$$2^\circ \text{ addend} = \int_{E(t)} \frac{\partial \ell}{\partial t} (z, t) dz$$

$$\text{Recall that : } G_t(x), G_t^{-1}(y) \quad \left(DG_t(G_t^{-1}(y)) \right)^{-1} = DG_t^{-1}(y)$$

$$3^\circ \text{ addend} =$$

$$\begin{aligned}
&= \int_{E_0} \rho (G_t(x), t) |DG_t(x)| \operatorname{tr} \left((DG_t(x))^{-1} \frac{\partial}{\partial t} DG_t(x) \right) dx =
\end{aligned}$$

$$\begin{aligned}
&= \int_{E_0} \rho (G_t(x), t) |DG_t(x)| \operatorname{tr} \left(DG_t^{-1}(G_t(x)) \cdot \frac{\partial}{\partial t} DG_t(x) \right) dx =
\end{aligned}$$

$$\text{Now } \left(D\zeta_t^{-1} \right)_{ij} (\zeta_t(x)) = \frac{\partial (\zeta_t^{-1})_i}{\partial y_j} (\zeta_t(x))$$

$$\frac{\partial}{\partial t} (D\zeta_t)_{jk}^{(x)} = \frac{\partial}{\partial t} \frac{\partial (\zeta_t)_j}{\partial x_k} (x, t)$$

$$\text{then } \left(D\zeta_t^{-1} (\zeta_t(x)) \cdot \frac{\partial}{\partial t} D\zeta_t^{(x)} \right)_{ik} =$$

$$= \sum_{j=1}^m \frac{\partial (\zeta_t^{-1})_i}{\partial y_j} (\zeta_t(x)) \frac{\partial^2 (\zeta_t)_j}{\partial t \partial x_k} (x)$$

$$(\text{for } i=k) = \sum_{j=1}^m \left(\frac{\partial}{\partial t} \frac{\partial (\zeta_t)_j}{\partial x_i} (x) \right) \frac{\partial (\zeta_t^{-1})_i}{\partial y_j} (\zeta_t(x))$$

$$= \sum_{j=1}^m \frac{\partial}{\partial x_i} \frac{\partial (\zeta_t)_j}{\partial t} (x) \frac{\partial (\zeta_t^{-1})_i}{\partial y_j} (\zeta_t(x))$$

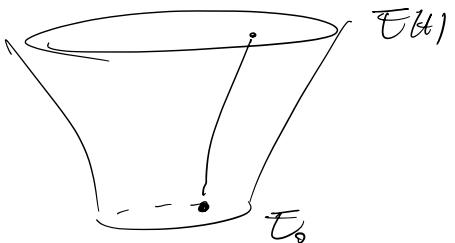
$$\text{tr} \left(D\zeta_t^{-1} (\zeta_t(x)) \cdot \frac{\partial}{\partial t} D\zeta_t^{(x)} \right) =$$

$$= \sum_{i=1}^n \sum_{j=1}^m \frac{\partial}{\partial x_i} \frac{\partial (\zeta_t)_j}{\partial t} (x, t) \frac{\partial (\zeta_t^{-1})_i}{\partial y_j} (\zeta_t(x))$$

Replacing $\zeta_t(x)$ with y we get

$$\sum_{j=1}^m \frac{\partial}{\partial y_j} \left(\frac{\partial (\zeta_t)_j}{\partial t} (\zeta_t^{-1}(y)) \right) = \text{div } V(y, t)$$

$$\text{where } \sqrt_j(y, t) := \frac{\partial(G_t)_j}{\partial t}(G_t^{-1}(y), t)$$



Then the 1st addend and
the 3rd addend give

$$\begin{aligned}
 & \int_{E_0} \sum_{i=1}^n \frac{\partial \rho}{\partial y_i} (G_t(x), t) \frac{\partial (G_t)_i}{\partial t} (x, t) |DG_t(x)| dx + \\
 & + \int_{E_0} \rho (G_t(x), t) |DG_t(x)| \operatorname{tr} \left((DG_t(x))^{-1} \cdot \frac{\partial}{\partial t} DG_t(x) \right) dx = \\
 & = \int_{E(t)} \sum_{i=1}^n \frac{\partial \rho}{\partial y_i} (y, t) \frac{\partial (G_t)_i}{\partial t} (y, t) dy + \\
 & + \int_{E(t)} \rho (y, t) \operatorname{div} \sqrt{\rho(y, t)} dy = \\
 & = \int_{E(t)} \operatorname{div} (\rho(y, t) \sqrt{\rho(y, t)}) dy
 \end{aligned}$$

$\sqrt{\rho}$ vector field denoting the velocity of $x_0 \in E_0$ at time t

Then we conclude that

$$\frac{d}{dt} \int_{E(t)} \rho(x,t) dx =$$

$$= \int_{E(t)} \frac{\partial \rho}{\partial t}(y,t) dy + \int_{E(t)} \operatorname{div}(\rho(y,t) \nabla(y,t)) dy$$

Now suppose ρ is the density of a fluid

($\rho(\cdot, t)$ density at time t)

$E(t)$ region occupied by the fluid at time t

examples: - ρ density of a compressible gas

- " " a liquid (water in the ground)
that is incompressible

Now whatever the region $E(t)$ the mass is
preserved, that is

$$\int_{E(t)} \rho(x, t) dx = \text{constant}$$

$E(t)$

then $\frac{d}{dt} \int_{E(t)} \rho(x, t) dx = 0$

and since
this holds
for every
choice of
 $E(t)$

We get
$$\frac{\partial \rho}{\partial t} + \operatorname{div}(\rho \mathbf{v}) = 0$$

This is called equation of continuity or of conservation of mass.

let us see some examples.

TRANSPORT EQUATION

$$n=1 \quad V = b \in \mathbb{R} \quad (\text{constant})$$

$$\frac{\partial \varphi}{\partial t} + b \frac{\partial \varphi}{\partial x} = 0$$

$$n > 1 \quad V \text{ constant} \quad (V = (V_1, \dots, V_n) \in \mathbb{R}^n)$$

$$\frac{\partial \varphi}{\partial t} + V \cdot \nabla \varphi = 0$$

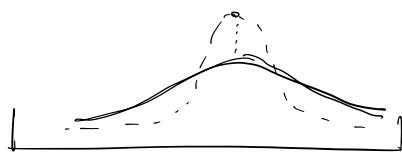
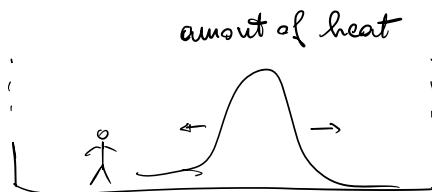
HEAT EQUATION

u Temperature

the Fourier's law

$$u_t = -k \nabla u$$

$k \in \mathbb{R}$
 $k > 0$
(conductivity)



$$\frac{\partial u}{\partial t} - k \operatorname{div}(\nabla u) =$$

$$\boxed{\frac{\partial u}{\partial t} - k \Delta u = 0}$$

LAPLACE EQUATION

ρ density of an incompressible fluid (\approx liquid)

Darcy's law $\mathbf{V} = -k \nabla p$ (fluid in a porous medium)

where p is
the pressure

$$\rho_t = 0 \quad \operatorname{div}(\rho k \nabla p) = \rho k \Delta p$$

$k > 0$

(because of
incompressibility,
 ρ constant)

$$-\Delta p = 0$$

WAVE EQUATION

transport in two directions, for instance

$$u_t + u_x = 0$$

$$u_t - u_x = 0$$

$$\left(\frac{\partial}{\partial t} + \frac{\partial}{\partial x} \right) u = 0 \quad \left(\frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u = 0$$

$$\boxed{u_{tt} - u_{xx} = 0}$$

can be factorized as

$$\left(\frac{\partial}{\partial t} + \frac{\partial}{\partial x} \right) \left(\frac{\partial}{\partial t} - \frac{\partial}{\partial x} \right) u = 0$$

A more general equation is

$$(k > 0, b \in \mathbb{R}^n, c \in \mathbb{R})$$

$$a \frac{\partial u}{\partial t} - k \Delta u + b \cdot \nabla u + cu = 0$$

diffusion advection
or transport reaction

THE NEXT THINGS WERE NOT DONE IN CLASS

THE PROOF IS GIVEN FOR THE SAKE OF COMPLETENESS

Proof of

$$\frac{d}{dt} (\det A(t)) = \det A(t) \cdot \text{tr} \left(A^{-1}(t) \frac{d}{dt} A(t) \right)$$

dim, consider the matrix

$$\det(\text{Id} + H) = \det(e_1 + H_1, \dots, e_n + H_n) \quad \begin{matrix} \text{Id} + H \\ \downarrow \\ \text{column vectors} \end{matrix}$$

because of
multilinearity

$$= \det(e_1, e_2 + H_2, e_3 + H_3, \dots) + \\ + \det(H_1, e_2 + H_2, e_3 + H_3, \dots) =$$

$$= \det(e_1, e_2, e_3 + H_3, \dots) + \det(e_1, H_2, e_3 + H_3, \dots) \\ + \det(H_1, e_2, e_3 + H_3, \dots) + \det(H_1, H_2, e_3 + H_3, \dots)$$

= ...

Observe that

$$\det(H_1, e_2, e_3, \dots, e_n) = h_{11}$$

$$\det(e_1, H_2, e_3, \dots, e_n) = h_{22}$$

:

$$\det(e_1, e_2, \dots, e_n, H_n) = h_{nn}$$

whose sum is the trace of H

$$\det(h_1, h_2, e_3, \dots e_n) = h_{11}h_{22} - h_{12}h_{21}$$

Then

$$\begin{aligned}\det(Id + H) &= \det Id + \text{tr } H + \\ &\quad + \text{ terms of order higher than 1} \\ &\quad \text{in } h_{ij}\end{aligned}$$

$$" = " 1 + \text{tr } H + o(H)$$

Recall : $f: \mathbb{R} \rightarrow \mathbb{R}$

$$\text{if } f(x) = f(x_0) + \alpha + o(|x-x_0|)$$

$$\text{then } \alpha = f'(x_0)(x-x_0)$$

$$F: \mathbb{R}^n \rightarrow \mathbb{R}$$

$$\text{if } F(x) = F(x_0) + \alpha + o(|x-x_0|)$$

$$\text{then } \alpha = (\nabla F(x_0), x-x_0)$$

$$F(x_0+v) = F(x_0) + \alpha + o(|x-x_0|)$$

$$\Rightarrow \alpha = \frac{\partial F}{\partial v}(x_0) = dF_{x_0} v$$

Then, if we denote by $\bar{F}(A)$ the function $\det A$,

$$\text{we have that } \frac{\partial \bar{F}}{\partial H}(\text{Id}) = d\bar{F}_{\text{Id}}H = \text{tr } H$$

$$\bar{F}(\text{Id} + H) = 1 + \text{tr } H + o(H)$$

$$\det(A + H) = \det(\lambda(\text{Id} + A^{-1}H)) =$$

$$= \det A \cdot \det(\text{Id} + A^{-1}H) =$$

$$= \det A \left(1 + \text{tr } A^{-1}H + \underset{o(H)}{\circ}(A^{-1}H) \right)$$

$$\bar{F}(A + H) = \det A + \det A \cdot \text{tr } A^{-1}H + o(H)$$

$$\frac{\partial \bar{F}}{\partial H}(A) = d\bar{F}_A(H) = \det A \quad \text{tr } A^{-1}H$$

$$\frac{d}{dt} \bar{F}(A(t)) = \det A(t) \quad \text{tr} \left(A(t)^{-1} \frac{d}{dt} A(t) \right) \quad \text{(*)}$$

$$\frac{d}{dt} (\det A(t))$$