# A higher dimensional Poincaré - Birkhoff theorem for Hamiltonian flows

Alessandro Fonda

(Università degli Studi di Trieste)

# A higher dimensional Poincaré - Birkhoff theorem for Hamiltonian flows

Alessandro Fonda

(Università degli Studi di Trieste)

a collaboration with Antonio J. Ureña

# A higher dimensional Poincaré - Birkhoff theorem for Hamiltonian flows

Alessandro Fonda

(Università degli Studi di Trieste)

a collaboration with Antonio J. Ureña

Annales de l'Institut Henri Poincaré (2017)

## But before starting...

But before starting...

### let me show you two recent photos...

# Oberwolfach, 1985



# Along the Adriatic, 1987



### Ok, let's start now

# Jules Henri Poincaré (1854 – 1912)



SUR UN THÉORÈME DE GÉOMÉTRIE.

Par M. H. Poincaré (Paris).

Adunanza del 10 marzo 1912.



Rend. Circ. Matem. Palermo, t. XXXIII (1º sem. 1912). - Stampato il 7 maggio 1912.



Rend. Circ. Matem. Palermo, t. XXXIII (1º sem. 1912). - Stampato il 7 maggio 1912.

Note: Poincaré died on July 17th, 1912

# RENDICONTI

DEL

# CIRCOLO MATEMATICO

### DI PALERMO

DIRETTORE: G. B. GUCCIA.

TOMO XXXIII (1° SEMESTRE 1912).

#### COMITATO DI REDAZIONE PEL TRIENNIO 1912-1913-1914.

#### Residenti:

M. L. Albeggiani. — G. Bagnera. — M. Gebbia. — G. B. Guccia. — G. Scorza.

#### Non Residenti:

E. BERTINI (Pisa). — L. BIANCHI (Pisa). — É. BOREL (Paris). — C. CARATHÉO-DORY (Breslau). — G. CASTELNUOVO (ROMA). — M. DE FRANCHIS (CAtania). — Ú. DINI (Pisa). — F. ENRIQUES (Bologna). — L. FEJÉR (Budapest). — A. R. FORSYTH (Cambridge). — I. FREDHOLM (Stockholm). — J. HADAMARD (Paris). — D. HILBERT (Göttingen). — G. HUMBERT (Paris). — F. KLEIN (Göttingen). — E. LANDAU (Göttingen). — T. LEVI-CIVITA (Padova). — A. LIAPOUNOFF (St.-Pétersburg). — G. LORIA (GENOVA). — A. E. H. LOVE (Oxford). — R. MARCOLONGO (Napoli). — F. MERTENS (Wien). — G. MITTAG-LEFFLER (Stockholm). — E. H. MOORE (Chicago, Ill.). — M. NOETHER (Erlangen). — W. F. OSGOOD (Cambridge, Mass.). — E. PASCAI (Napoli). — É. PICARD (Paris). — S. PINCHERLE (Bologna). — H. POINCARÉ (Paris). — C. SEGRE (TORINO). — F. SEVERI (Padova). — C. SOMIGLIANA (TORINO). — P. STÄCKEL (KARISTUR). — W. STEKLOFF (St.-Pétersbourg). — C. STÉPHANOS (Athènes). — Ch.-J. DE LA VALLÉE POUSSIN (Louvain). — G. VIVANTI (Pavia). — W. WIRTINGER (Wien). — H. G. ZEUTHEN (Kóbenhavn).

Direttore dei Rendiconti: G. B. GUCCIA.

### COMITATO DI REDAZIONE

PEL TRIENNIO 1912-1913-1914.

#### Residenti:

M. L. Albeggiani. — G. Bagnera. — M. Gebbia. — G. B. Guccia. — G. Scorza.

Non Residenti:

E. BERTINI (Pisa). — L. BIANCHI (Pisa). — É. BOREL (Paris). — C. CARATHÉO-DORY Breslau). — G. CASTELNUOVO (Roma). — M. DE FRANCHIS (Catania). — U. DINI (Pisa).—F. ENRIQUES (Bologna). — L. FEJÉR (Budapest). — A. R. FORSYTH (Cambridge). — I. FREDHOLM (Stockholm). — J. HADAMARD (Paris). — D. HILBERT (Göttingen). — G. HUMBERT (Paris). — F. KLEIN (Göttingen).—E. LANDAU (Göttingen). — T. LEVI-CIVITA (Padova). — A. LIAPOUNOFF (St.-Pétersbourg). — G. LORIA (Genova). — A. E. H. LOVE (Oxford). — R. MARCOLONGO (Napoli). — F. MERTENS (Wien). — G. MITTAG-LEFFLER (Stockholm).—E. H. MOORE (Chicago, Ill.).—M. NOETHER (Erlangen).—W. F. OSGOOD (Cambridge, Mass.).—E. PASCAL (Napoli). — E. PICARD (Paris).—S. PINCHERLE (Bologna). — H. POINCARÉ (Paris).—C. SEGRE (Torino).—F. SEVERI (Padova).—C. SOMIGLIANA (Torino).—P. STÄCKEL (KARISTUR).—W. STEKLOFF (St.-Pétersbourg).—C. STÉPHANOS (Athènes). — Ch.-J. DE LA VALLÉE POUSSIN (Louvain). — G. VIVANTI (Pavia). — W. WIRTINGER (Wien). — H. G. ZEUTHEN (Kobenhavn).

Direttore dei Rendiconti: G. B. GUCCIA.

#### SUR UN THÉORÈME DE GÉOMÉTRIE.

Par M. H. Poincaré (Paris).

Adunanza del 10 marzo 1912.

#### § I.

#### INTRODUCTION.

Je n'ai jamais présenté au public un travail aussi inachevé; je crois donc nécessaire d'expliquer en quelques mots les raisons qui m'ont déterminé à le publier, et d'abord celles qui m'avaient engagé à l'entreprendre. J'ai démontré, il y a longtemps déjà, l'existence des solutions périodiques du problème des trois corps; le résultat laissait cependant encore à désirer; car, si l'existence de chaque sorte de solution était établie pour les petites valeurs des masses, on ne voyait pas ce qui devait arriver pour des valeurs plus grandes, quelles étaient celles de ces solutions qui subsistaient et dans quel ordre elles disparaissaient. En réfléchissant à cette question, je me suis assuré que la réponse devait dépendre de l'exactitude ou de la fausseté d'un certain théorème de géométrie dont l'énoncé est très simple, du moins dans le cas du problème restreint et des problèmes de Dynamique où il n'y a que deux degrés de liberté.

#### SUR UN THÉORÈME DE GÉOMÉTRIE.

Par M. H. Poincaré (Paris).

Adunanza del 10 marzo 1912.

#### § I.

#### INTRODUCTION.

Je n'ai jamais présenté au public un travail aussi inachevé je crois donc nécessaire d'expliquer en quelques mots les raisons qui m'ont déterminé à le publier, et d'abord celles qui m'avaient engagé à l'entreprendre. l'ai démontré, il y a longtemps déjà, l'existence des solutions périodiques du problème des trois corps le résultat laissait cependant encore à désirer; car, si l'existence de chaque sorte de solution était établie pour les petites valeurs des masses, on ne voyait pas ce qui devait arriver pour des valeurs plus grandes, quelles étaient celles de ces solutions qui subsistaient et dans que ordre elles disparaissaient. En réfléchissant à cette question, je me suis assuré que la réponse devait dépendre de l'exactitude ou de la fausseté d'un certain théorème de géométrie dont l'énoncé est très simple du moins dans le cas du problème restreint et des problèmes de Dynamique où il n'y a que deux degrés de liberté. J'ai donc été amené à rechercher si ce théorème est vrai ou faux, mais j'ai rencontré des difficultés auxquelles je ne m'attendais pas. J'ai été obligé d'envisager séparément un très grand nombre de cas particuliers; mais les cas possibles sont trop nombreux pour que j'aie pu les étudier tous. J'ai reconnu l'exactitude du théorème dans tous ceux que j'ai traités. Pendant deux ans, je me suis efforcé sans succès, soit de trouver une démonstration générale, soit de découvrir un exemple où le théorème soit en défaut.

Ma conviction qu'il est toujours vrai s'affermissait de jour en jour, mais je restais incapable de l'asseoir sur des fondements solides.

Il semble que dans ces conditions, je devrais m'abstenir de toute publication tant que je n'aurai pas résolu la question; mais après les inutiles efforts que j'ai faits pendant de longs mois, il m'a paru que le plus sage était de laisser le problème mûrir, en m'en reposant durant quelques années; cela serait très bien si j'étais sûr de pouvoir le reprendre un jour; mais à mon âge je ne puis en répondre. D'un autre côté, l'importance du sujet est trop grande (et je chercherai plus loin à la faire comprendre) et l'ensemble des résultats obtenus trop considérable déjà, pour que je me résigne à les laisser définitivement infructueux. Je puis espérer que les géomètres qui s'intéresseront à ce pro-

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

and

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

and

 $(\star)$  it rotates the two boundary circles in opposite directions

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

and

(\*) it rotates the two boundary circles in opposite directions (this is called the "twist condition").

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

and

(\*) it rotates the two boundary circles in opposite directions (this is called the "twist condition").

 $\ensuremath{\mathcal{A}}$  is a closed planar annulus



 $\mathcal{P}:\mathcal{A}\rightarrow\mathcal{A}$  is an area preserving homeomorphism

#### and

- (\*) it rotates the two boundary circles in opposite directions (this is called the "twist condition").
- Then,  $\mathcal{P}$  has two fixed points.

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S}$  is an area preserving homeomorphism

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S}$  is an area preserving homeomorphism, and writing

$$\mathcal{P}(x,y) = (x + f(x,y), y + g(x,y)),$$

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S} \, \text{ is an area preserving homeomorphism,} \\ and writing$ 

$$\mathcal{P}(x, y) = (x + f(x, y), y + g(x, y)),$$

both f(x, y) and g(x, y) are continuous,  $2\pi$ -periodic in x,

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S} \, \text{ is an area preserving homeomorphism,} \\ and writing$ 

$$\mathcal{P}(x, y) = (x + f(x, y), y + g(x, y)),$$

both f(x, y) and g(x, y) are continuous,  $2\pi$ -periodic in x,

g(x, a) = 0 = g(x, b) (boundary invariance),

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S} \, \text{ is an area preserving homeomorphism,} \\ and writing$ 

$$\mathcal{P}(x, y) = (x + f(x, y), y + g(x, y)),$$

both f(x, y) and g(x, y) are continuous,  $2\pi$ -periodic in x,

g(x, a) = 0 = g(x, b) (boundary invariance),

and

(\*) 
$$f(x, a) < 0 < f(x, b)$$
 (twist condition).

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S} \, \text{ is an area preserving homeomorphism,} \\ and writing$ 

$$\mathcal{P}(x, y) = (x + f(x, y), y + g(x, y)),$$

both f(x, y) and g(x, y) are continuous,  $2\pi$ -periodic in x,

g(x, a) = 0 = g(x, b) (boundary invariance),

and

(\*) 
$$f(x, a) < 0 < f(x, b)$$
 (twist condition).

 $\mathcal{S} = \mathbb{R} \times [a, b]$  is a planar strip



 $\mathcal{P}: \mathcal{S} \rightarrow \mathcal{S}$  is an area preserving homeomorphism, and writing

$$\mathcal{P}(x, y) = (x + f(x, y), y + g(x, y)),$$

both f(x, y) and g(x, y) are continuous,  $2\pi$ -periodic in x,

g(x, a) = 0 = g(x, b) (boundary invariance),

and

(\*) 
$$f(x, a) < 0 < f(x, b)$$
 (twist condition).

Then,  $\mathcal{P}$  has two geometrically distinct fixed points.

# George David Birkhoff (1884 - 1944)


# The Poincaré – Birkhoff theorem

In 1913 – 1925, Birkhoff proved Poincaré's "théorème de géométrie", so that it now carries the name

"Poincaré – Birkhoff Theorem".

# The Poincaré – Birkhoff theorem

In 1913 – 1925, Birkhoff proved Poincaré's "théorème de géométrie", so that it now carries the name

"Poincaré – Birkhoff Theorem".

Variants and different proofs have been proposed by:

Brown–Neumann, Carter, W.-Y. Ding, Franks, Guillou, Jacobowitz, de Kérékjartó, Le Calvez, Moser, Rebelo, Slaminka, ...

# The Poincaré – Birkhoff theorem

In 1913 – 1925, Birkhoff proved Poincaré's "théorème de géométrie", so that it now carries the name

"Poincaré – Birkhoff Theorem".

Variants and different proofs have been proposed by:

Brown–Neumann, Carter, W.-Y. Ding, Franks, Guillou, Jacobowitz, de Kérékjartó, Le Calvez, Moser, Rebelo, Slaminka, ...

Applications to the existence of periodic solutions were provided by:

Bonheure, Boscaggin, Butler, Del Pino, T. Ding, Fabry, Garrione, Hartman, Manásevich, Mawhin, Omari, Sfecci, Smets, Torres, Wang, Zanini, Zanolin, ...

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is *T*-periodic in *t*. The questions we want to face:

Are there periodic solutions? How many?

We consider the system

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{y}}(t, \mathbf{x}, \mathbf{y}), \qquad \dot{\mathbf{y}} = -\frac{\partial H}{\partial \mathbf{x}}(t, \mathbf{x}, \mathbf{y}),$$

and assume that the Hamiltonian H(t, x, y) is *T*-periodic in *t*. The questions we want to face:

Are there periodic solutions? How many?

Two "simple" examples: the pendulum equation

$$\ddot{x}+\sin x=e(t),$$

and the superlinear equation

$$\ddot{x} + x^3 = e(t) \,,$$

where e(t) is a *T*-periodic forcing.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

The Poincaré time - map is defined as

$$\mathcal{P}: (x_0, y_0) \mapsto (x_T, y_T)$$

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

The Poincaré time - map is defined as

$$\mathcal{P}: (x_0, y_0) \mapsto (x_T, y_T)$$

i.e.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

The Poincaré time - map is defined as

$$\mathcal{P}: (x_0, y_0) \mapsto (x_T, y_T)$$

i.e.

to each "starting point"  $(x_0, y_0)$  of a solution at time t = 0,

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

The Poincaré time - map is defined as

$$\mathcal{P}: (x_0, y_0) \mapsto (x_T, y_T)$$

#### i.e.

to each "starting point"  $(x_0, y_0)$  of a solution at time t = 0,

 $\ensuremath{\mathcal{P}}$  associates

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

The Poincaré time - map is defined as

$$\mathcal{P}: (x_0, y_0) \mapsto (x_T, y_T)$$

i.e.

to each "starting point"  $(x_0, y_0)$  of a solution at time t = 0,

 $\ensuremath{\mathcal{P}}$  associates

the "arrival point"  $(x_T, y_T)$  of the solution at time t = T.

## Good and bad news

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

## Good and bad news

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Good news:

The Poincaré map  $\mathcal{P}$  is an area preserving homeomorphism.

Its fixed points correspond to *T*-periodic solutions.

## Good and bad news

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

#### Good news:

The Poincaré map  $\ensuremath{\mathcal{P}}$  is an area preserving homeomorphism.

Its fixed points correspond to *T*-periodic solutions.

Bad news:

It is very difficult to find an invariant annulus for  $\ensuremath{\mathcal{P}}$  .

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

We consider the system

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{y}}(t, \mathbf{x}, \mathbf{y}), \qquad \dot{\mathbf{y}} = -\frac{\partial H}{\partial \mathbf{x}}(t, \mathbf{x}, \mathbf{y}),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Assume H(t, x, y) to be also  $2\pi$ -periodic in x.

We consider the system

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{y}}(t, \mathbf{x}, \mathbf{y}), \qquad \dot{\mathbf{y}} = -\frac{\partial H}{\partial \mathbf{x}}(t, \mathbf{x}, \mathbf{y}),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Assume H(t, x, y) to be also  $2\pi$ -periodic in x.

Let  $S = \mathbb{R} \times [a, b]$  be a planar strip.

We consider the system

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{y}}(t, \mathbf{x}, \mathbf{y}), \qquad \dot{\mathbf{y}} = -\frac{\partial H}{\partial \mathbf{x}}(t, \mathbf{x}, \mathbf{y}),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Assume H(t, x, y) to be also  $2\pi$ -periodic in x.

Let  $S = \mathbb{R} \times [a, b]$  be a planar strip.

Twist condition: the solutions (x(t), y(t)) with "starting point" (x(0), y(0)) on  $\partial S$  are defined on [0, T] and satisfy

(\*) 
$$x(T) - x(0) \begin{cases} < 0, & \text{if } y(0) = a, \\ > 0, & \text{if } y(0) = b. \end{cases}$$

We consider the system

$$\dot{\mathbf{x}} = \frac{\partial H}{\partial \mathbf{y}}(t, \mathbf{x}, \mathbf{y}), \qquad \dot{\mathbf{y}} = -\frac{\partial H}{\partial \mathbf{x}}(t, \mathbf{x}, \mathbf{y}),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Assume H(t, x, y) to be also  $2\pi$ -periodic in x.

Let  $\mathcal{S} = \mathbb{R} \times [a, b]$  be a planar strip.

Twist condition: the solutions (x(t), y(t)) with "starting point" (x(0), y(0)) on  $\partial S$  are defined on [0, T] and satisfy

(\*) 
$$x(T) - x(0) \begin{cases} < 0, & \text{if } y(0) = a, \\ > 0, & \text{if } y(0) = b. \end{cases}$$

Then, there are two geometrically distinct T-periodic solutions.

1. Writing  $\mathcal{S} = \mathbb{R} \times \overline{\mathcal{D}}$ , with

 $\mathcal{D} = ]a, b[,$ 

1. Writing  $\mathcal{S} = \mathbb{R} \times \overline{\mathcal{D}}$ , with

$$\mathcal{D} = ]a, b[,$$

and defining the "outer normal function"  $\nu: \partial \mathcal{D} \to \mathbb{R}$  as

$$u(a) = -1, \quad \nu(b) = +1,$$

1. Writing 
$$\mathcal{S} = \mathbb{R} \times \overline{\mathcal{D}}$$
, with

$$\mathcal{D} = ]a, b[,$$

and defining the "outer normal function"  $\nu: \partial \mathcal{D} \to \mathbb{R}$  as

$$u(a) = -1, \quad \nu(b) = +1,$$

the twist condition

$$(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad x(T) - x(0) \begin{cases} < 0, & \text{if } y(0) = a, \\ > 0, & \text{if } y(0) = b, \end{cases}$$

1. Writing 
$$\mathcal{S} = \mathbb{R} \times \overline{\mathcal{D}}$$
, with

$$\mathcal{D} = ]a, b[,$$

and defining the "outer normal function"  $\nu: \partial \mathcal{D} \to \mathbb{R}$  as

$$u(a) = -1, \quad \nu(b) = +1,$$

the twist condition

$$(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad x(T) - x(0) \begin{cases} < 0, & \text{if } y(0) = a, \\ > 0, & \text{if } y(0) = b, \end{cases}$$

can be written as

$$(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad [x(T) - x(0)] \cdot \nu(y(0)) > 0.$$

1. Writing 
$$\mathcal{S} = \mathbb{R} \times \overline{\mathcal{D}}$$
, with

$$\mathcal{D} = ]a, b[,$$

and defining the "outer normal function"  $\nu: \partial \mathcal{D} \to \mathbb{R}$  as

$$u(a) = -1, \quad \nu(b) = +1,$$

the twist condition

$$(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad x(T) - x(0) \begin{cases} < 0, & \text{if } y(0) = a, \\ > 0, & \text{if } y(0) = b, \end{cases}$$

can be written as

$$(\star) \qquad (x(0), y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad [x(\mathcal{T}) - x(0)] \cdot \nu(y(0)) > 0 \,.$$

2. The Poincaré map could be multivalued.

The outstanding question as to the possibility of an N-dimensional extension of Poincaré's last geometric theorem

[Birkhoff, Acta Mathematica 1925]

The outstanding question as to the possibility of an N-dimensional extension of Poincaré's last geometric theorem

[Birkhoff, Acta Mathematica 1925]

Attempts in some directions have been made by:

Amann, Bertotti, Birkhoff, K.C. Chang, Conley, Felmer, Golé, Hingston, Josellis, J.Q. Liu, Mawhin, Moser, Rabinowitz, Szulkin, Weinstein, Willem, Winkelnkemper, Zehnder, ...

The outstanding question as to the possibility of an N-dimensional extension of Poincaré's last geometric theorem

[Birkhoff, Acta Mathematica 1925]

Attempts in some directions have been made by:

Amann, Bertotti, Birkhoff, K.C. Chang, Conley, Felmer, Golé, Hingston, Josellis, J.Q. Liu, Mawhin, Moser, Rabinowitz, Szulkin, Weinstein, Willem, Winkelnkemper, Zehnder, ...

However,

a genuine generalization of the Poincaré – Birkhoff theorem to higher dimensions has never been given.

[Moser and Zehnder, Notes on Dynamical Systems, 2005].

The outstanding question as to the possibility of an N-dimensional extension of Poincaré's last geometric theorem

[Birkhoff, Acta Mathematica 1925]

Attempts in some directions have been made by:

Amann, Bertotti, Birkhoff, K.C. Chang, Conley, Felmer, Golé, Hingston, Josellis, J.Q. Liu, Mawhin, Moser, Rabinowitz, Szulkin, Weinstein, Willem, Winkelnkemper, Zehnder, ...

However,

a genuine generalization of the Poincaré – Birkhoff theorem to higher dimensions has never been given.

[Moser and Zehnder, Notes on Dynamical Systems, 2005].

Note: Arnold proposed some conjectures in the sixties. Some of them are still open.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here,  $x = (x_1, ..., x_N)$  and  $y = (y_1, ..., y_N)$ .

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here, 
$$x = (x_1, ..., x_N)$$
 and  $y = (y_1, ..., y_N)$ .

Assume H(t, x, y) to be also  $2\pi$ -periodic in each  $x_1, \ldots, x_N$ .
We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here, 
$$x = (x_1, ..., x_N)$$
 and  $y = (y_1, ..., y_N)$ .

Assume H(t, x, y) to be also  $2\pi$ -periodic in each  $x_1, \ldots, x_N$ .

Let  $\mathcal{D}$  be an open, bounded, convex set in  $\mathbb{R}^N$ , with a smooth boundary, and denote by  $\nu : \partial \mathcal{D} \to \mathbb{R}^N$  the outward normal vectorfield. Consider the "strip"  $\mathcal{S} = \mathbb{R}^N \times \overline{\mathcal{D}}$ .

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here, 
$$x = (x_1, ..., x_N)$$
 and  $y = (y_1, ..., y_N)$ .

Assume H(t, x, y) to be also  $2\pi$ -periodic in each  $x_1, \ldots, x_N$ .

Let  $\mathcal{D}$  be an open, bounded, convex set in  $\mathbb{R}^N$ , with a smooth boundary, and denote by  $\nu : \partial \mathcal{D} \to \mathbb{R}^N$  the outward normal vectorfield. Consider the "strip"  $\mathcal{S} = \mathbb{R}^N \times \overline{\mathcal{D}}$ .

Twist condition: for a solution (x(t), y(t)),

$$(\star) \qquad (x(0),y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad [x(\mathcal{T})-x(0)] \cdot \nu(y(0)) > 0 \,.$$

(this is the old condition, when N = 1)

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here, 
$$x = (x_1, ..., x_N)$$
 and  $y = (y_1, ..., y_N)$ .

Assume H(t, x, y) to be also  $2\pi$ -periodic in each  $x_1, \ldots, x_N$ .

Let  $\mathcal{D}$  be an open, bounded, convex set in  $\mathbb{R}^N$ , with a smooth boundary, and denote by  $\nu : \partial \mathcal{D} \to \mathbb{R}^N$  the outward normal vectorfield. Consider the "strip"  $\mathcal{S} = \mathbb{R}^N \times \overline{\mathcal{D}}$ .

Twist condition: for a solution (x(t), y(t)),

$$(\star) \qquad (x(0),y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad \left\langle x(T) - x(0), \nu(y(0)) \right\rangle > 0 \,.$$

(this is the new condition)

We consider the system

$$\dot{x} = \frac{\partial H}{\partial y}(t, x, y), \qquad \dot{y} = -\frac{\partial H}{\partial x}(t, x, y),$$

and assume that the Hamiltonian H(t, x, y) is T-periodic in t.

Here, 
$$x = (x_1, ..., x_N)$$
 and  $y = (y_1, ..., y_N)$ .

Assume H(t, x, y) to be also  $2\pi$ -periodic in each  $x_1, \ldots, x_N$ .

Let  $\mathcal{D}$  be an open, bounded, convex set in  $\mathbb{R}^N$ , with a smooth boundary, and denote by  $\nu : \partial \mathcal{D} \to \mathbb{R}^N$  the outward normal vectorfield. Consider the "strip"  $\mathcal{S} = \mathbb{R}^N \times \overline{\mathcal{D}}$ .

Twist condition: for a solution (x(t), y(t)),

$$(\star) \qquad (x(0),y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad \left\langle x(T) - x(0), \nu(y(0)) \right\rangle > 0 \,.$$

Then, there are N + 1 geometrically distinct T-periodic solutions.

The proof is variational, it uses an

infinite dimensional Ljusternik – Schnirelmann theory.

The proof is variational, it uses an

infinite dimensional Ljusternik – Schnirelmann theory.

The periodicity in  $x_1, ..., x_N$  permits to define the action functional on the product of a Hilbert space *E* and the *N*-torus  $\mathbb{T}^N$ :

 $\varphi: \boldsymbol{E} \times \mathbb{T}^{\boldsymbol{N}} \to \mathbb{R} \, .$ 

The proof is variational, it uses an

infinite dimensional Ljusternik - Schnirelmann theory.

The periodicity in  $x_1, ..., x_N$  permits to define the action functional on the product of a Hilbert space *E* and the *N*-torus  $\mathbb{T}^N$ :

 $\varphi: \boldsymbol{E} \times \mathbb{T}^{\boldsymbol{N}} \to \mathbb{R} \, .$ 

The result then follows from the fact that

 $\operatorname{cat}(\mathbb{T}^N) = N + 1$ .

The proof is variational, it uses an

infinite dimensional Ljusternik – Schnirelmann theory.

The periodicity in  $x_1, \ldots, x_N$  permits to define the action functional on the product of a Hilbert space *E* and the *N*-torus  $\mathbb{T}^N$ :

 $\varphi: \boldsymbol{E} \times \mathbb{T}^N \to \mathbb{R}$ .

The result then follows from the fact that

$$\operatorname{cat}(\mathbb{T}^N) = N + 1$$
.

Note. If  $\varphi$  only has nondegenerate critical points, then we can use Morse theory and find  $2^N$  solutions.

The proof is variational, it uses an

infinite dimensional Ljusternik – Schnirelmann theory.

The periodicity in  $x_1, \ldots, x_N$  permits to define the action functional on the product of a Hilbert space *E* and the *N*-torus  $\mathbb{T}^N$ :

 $\varphi: \boldsymbol{E} \times \mathbb{T}^N \to \mathbb{R}$ .

The result then follows from the fact that

$$\operatorname{cat}(\mathbb{T}^N) = N + 1$$
.

Note. If  $\varphi$  only has nondegenerate critical points, then we can use Morse theory and find  $2^N$  solutions. Indeed,

$$\operatorname{sb}(\mathbb{T}^N) = 2^N$$
.

The twist condition

$$(\star) \qquad (x(0), y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad \big\langle x(T) - x(0), \nu(y(0)) \big\rangle > 0$$

can be improved in two directions.

The twist condition

 $(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad \left\langle x(T) - x(0), \nu(y(0)) \right\rangle > 0$ 

can be improved in two directions.

I. The "indefinite twist" condition: for a regular symmetric  $N \times N$  matrix  $\mathbb{A}$ ,

$$(\star') \qquad (x(0), y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad \big\langle x(T) - x(0), \mathbb{A}\nu(y(0)) \big\rangle > 0 \,.$$

The twist condition

 $(\star) \qquad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad \left\langle x(T) - x(0), \nu(y(0)) \right\rangle > 0$ 

can be improved in two directions.

I. The "indefinite twist" condition: for a regular symmetric  $N \times N$  matrix  $\mathbb{A}$ ,

$$(\star') \qquad (x(0), y(0)) \in \partial \mathcal{S} \quad \Rightarrow \quad \big\langle x(T) - x(0), \mathbb{A}\nu(y(0)) \big\rangle > 0 \,.$$

II. The "avoiding rays" condition:

$$(\star'') \quad (x(0), y(0)) \in \partial S \quad \Rightarrow \quad x(T) - x(0) \notin \{-\lambda \nu(y(0)) : \lambda \ge 0\}.$$

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

A. Fonda, M. Garrione and P. Gidoni, Periodic perturbations of Hamiltonian systems, Advances in Nonlinear Analysis (2016)

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

- A. Fonda, M. Garrione and P. Gidoni, Periodic perturbations of Hamiltonian systems, Advances in Nonlinear Analysis (2016)
- A. Fonda and A. Sfecci,

Multiple periodic solutions of Hamiltonian systems confined in a box, Discrete and Continuous Dynamical Systems (2017)

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

A. Fonda, M. Garrione and P. Gidoni, Periodic perturbations of Hamiltonian systems, Advances in Nonlinear Analysis (2016)

A. Fonda and A. Sfecci,

Multiple periodic solutions of Hamiltonian systems confined in a box, Discrete and Continuous Dynamical Systems (2017)

A. Fonda and P. Gidoni,

An avoiding cones condition for the Poincaré–Birkhoff theorem, Journal of Differential Equations (2017)

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

A. Fonda, M. Garrione and P. Gidoni, Periodic perturbations of Hamiltonian systems, Advances in Nonlinear Analysis (2016)

A. Fonda and A. Sfecci,

Multiple periodic solutions of Hamiltonian systems confined in a box, Discrete and Continuous Dynamical Systems (2017)

A. Fonda and P. Gidoni,

An avoiding cones condition for the Poincaré–Birkhoff theorem, Journal of Differential Equations (2017)

A. Fonda and R. Toader,

Subharmonic solutions of Hamiltonian systems displaying some kind of sublinear growth, Advances in Nonlinear Analysis (2017)

A. Fonda and A. Sfecci,

Periodic solutions of weakly coupled superlinear systems, Journal of Differential Equations (2016)

A. Fonda, M. Garrione and P. Gidoni, Periodic perturbations of Hamiltonian systems, Advances in Nonlinear Analysis (2016)

A. Fonda and A. Sfecci,

Multiple periodic solutions of Hamiltonian systems confined in a box, Discrete and Continuous Dynamical Systems (2017)

A. Fonda and P. Gidoni,

An avoiding cones condition for the Poincaré–Birkhoff theorem, Journal of Differential Equations (2017)

A. Fonda and R. Toader,

Subharmonic solutions of Hamiltonian systems displaying some kind of sublinear growth, Advances in Nonlinear Analysis (2017)

#### A. Boscaggin, A. Fonda and M. Garrione,

An infinite-dimensional version of the Poincaré–Birkhoff theorem on the Hilbert cube, preprint 2017



# Buon compleanno!!!