On the spectral behavior of a biharmonic Steklov problem

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Let Ω be a domain in \mathbb{R}^N . We start by recalling the classical Steklov eigenvalue problem

$$\begin{cases} \Delta u = 0, & \text{in } \Omega, \\ \frac{\partial u}{\partial v} = \lambda u, & \text{on } \partial \Omega. \end{cases}$$

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If Ω is a bounded domain with Lipschitz boundary, then the spectrum is discrete and the eigenvalues form a sequence

$$0 = \lambda_1 < \lambda_2 \leq \cdots \leq \lambda_j \leq \cdots \nearrow +\infty.$$



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- Shape optimization 2: Weinstock's inequality states that "among all simply connected planar domains with fixed perimeter, λ₂ is maximized by a disk". The general case is an open problem.



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The Steklov problem can be formulated in a more general setting:

- Ω is a compact Riemannian manifold of dimension N ≥ 2 with boundary ∂Ω;
- Δ is the Laplace-Beltrami operator on functions on Ω ;
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A. Girouard and I. Polterovich. Spectral geometry of the Steklov problem. *arXiv:1411.6567*, 2014.

I. Chavel. Eigenvalues in Riemannian geometry. Academic Press, INC., 1984.



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What about **Steklov boundary conditions** for the biharmonic operator?





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- If Ω is the unit ball in \mathbb{R}^N and $\lambda = N 1$, we have the hinged plate problem for the associated Poisson problem $\Delta^2 u = f$.



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- If Ω is the unit ball in \mathbb{R}^N and $\lambda = N 1$, we have the hinged plate problem for the associated Poisson problem $\Delta^2 u = f$.
- Shape optimization: the behavior is completely different from the second order case and not well understood.



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- Consider the issue of the stability of the optimal shapes and quantify it.



Recall (also from the original Steklov paper) that "The eigenvalues of the Laplace operator with Steklov boundary conditions represent the squares of the natural frequencies of vibration of a membrane with a free frame and the mass of which is displaced uniformly on the boundary".



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Then we start from the model of a free vibrating plate with mass density $\rho > 0$ on Ω , i.e., a Neumann problem for the biharmonic operator with mass density ρ .



Let Ω be a bounded domain of class C^1 . Let $\tau \ge 0$ be a fixed constant and $\rho : \Omega \to \mathbb{R}^+$ a positive function. We consider the Neumann problem

$$\begin{cases} \Delta^2 u - \tau \Delta u = \lambda \rho u, & \text{in } \Omega, \\ \frac{\partial^2 u}{\partial \nu^2} = \tau \frac{\partial u}{\partial \nu} - \operatorname{div}_{\partial \Omega} \left(D^2 u \cdot \nu \right) - \frac{\partial \Delta u}{\partial \nu} = 0, & \text{on } \partial \Omega, \end{cases}$$





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We consider the Neumann problem with density ρ_{ε}

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For each fixed $\boldsymbol{\varepsilon}$ we have an increasing sequence

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Theorem

For all $j \in \mathbb{N}$, $\lim_{\varepsilon \to 0} \lambda_j(\varepsilon) = \lambda_j$, where λ_j are the eigenvalues of

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The proof of the convergence of the eigenvalues consists in showing that the resolvent operators of the Neumann problems with ρ_{ε} compactly converge to the resolvent operator of the limiting problem.



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Is this map Continuous? Differentiable? Analytic?



$$\begin{cases} \Delta^2 u - \tau \Delta u = 0, & \text{in } \Omega, \\ \frac{\partial^2 u}{\partial \nu^2} = 0, & \text{on } \partial \Omega, \\ \tau \frac{\partial u}{\partial \nu} - \operatorname{div}_{\partial \Omega} \left(D^2 u \cdot \nu \right) - \frac{\partial \Delta u}{\partial \nu} = \lambda u, & \text{on } \partial \Omega. \end{cases}$$
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Critical points?







The set of domains has not a linear structure, so what does it means differentiability?

Let Ω be a fixed domain of class C^1 and let

$$\Phi(\Omega) = \left\{ \phi \in \left(C^2(\overline{\Omega}) \right)^N : \phi \text{ injective and } \inf_{\Omega} |\det D\phi| > 0 \right\}.$$





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$$\phi\mapsto \lambda_j[\phi].$$

The space $\Phi(\Omega)$ is a linear space, so we can make Differential Calculus on it.



 $\mathcal{A}_{\Omega}[F] = \left\{ \phi \in \Phi(\Omega) : \lambda_{l}[\phi] \neq \lambda_{j}[\phi] \quad \forall j \in F, \ \forall l \in \mathbb{N} \setminus F \right\}$





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For example, if $F = \{1\}$, then $\mathcal{A}_{\Omega}[F] = \{\phi \in \Phi(\Omega) : \lambda_1[\phi] \text{ is simple}\}.$





$$\mathcal{R}_{\Omega}[F] = \left\{ \phi \in \Phi(\Omega) : \lambda_{l}[\phi] \neq \lambda_{j}[\phi] \quad \forall j \in F, \ \forall l \in \mathbb{N} \setminus F \right\}$$

For example, if $F = \{1\}$, then $\mathcal{A}_{\Omega}[F] = \{\phi \in \Phi(\Omega) : \lambda_1[\phi] \text{ is simple}\}.$

Then we consider the symmetric functions of the eigenvalues, for $s \in \{1, ..., |F|\}$

$$\Lambda_{F,s}[\phi] = \sum_{j_1 < \cdots < j_s \in F} \lambda_{j_1}[\phi] \cdots \lambda_{j_s}[\phi]$$





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Such functions turn out to be important objects of study in shape optimization problems.







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Why the symmetric functions of the eigenvalues? Example:

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Symmetric functions of $\lambda_1[\alpha_1, \alpha_2], \lambda_2[\alpha_1, \alpha_2]$ are even analytic.



Theorem (Analyticity)

Let Ω be a bounded domain in \mathbb{R}^N of class C^1 . Let F be a finite non-empty subset of \mathbb{N} . Then

i) The set $\mathcal{A}_{\Omega}[F]$ is open in $\Phi(\Omega)$.




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- i) The set $\mathcal{A}_{\Omega}[F]$ is open in $\Phi(\Omega)$.
- ii) The function $\Lambda_{F,s}[\phi]$ from $\mathcal{A}_{\Omega}[F]$ to \mathbb{R} is real analytic.





Theorem (Derivatives)

Let $\tilde{\phi} \in \mathcal{A}_{\Omega}[F]$ be such that $\lambda_j[\tilde{\phi}] = \lambda_F[\tilde{\phi}]$ for all $j \in F$ and such that $\tilde{\phi}(\Omega)$ is of class C^4 . Let $v_1, ..., v_{|F|}$ be a orthonormal basis of the eigenspace associated with $\lambda_F[\tilde{\phi}]$. Then

$$\begin{aligned} d|_{\phi=\tilde{\phi}}\left(\Lambda_{F,s}\right)\left[\psi\right] &= -\lambda_{F}^{s-1}[\tilde{\phi}] \binom{|F|-1}{s-1} \sum_{j=1}^{|F|} \int_{\partial \tilde{\phi}(\Omega)} \left(\lambda_{F}[\tilde{\phi}] K v_{j}^{2} \right. \\ &\left. + \lambda_{F}[\tilde{\phi}] \frac{\partial(v_{j}^{2})}{\partial v} - \tau |\nabla v_{j}|^{2} - |D^{2} v_{j}|^{2} \right) \psi \circ \tilde{\phi}^{(-1)} \cdot v d\sigma, \end{aligned}$$

for all $\psi \in (C^2(\overline{\Omega}))^N$, where K denotes the mean curvature of $\partial \tilde{\phi}(\Omega)$.



$$\min_{\mathcal{V}(\phi)=\text{const.}} \Lambda_{F,s}[\phi] \quad \text{or} \quad \max_{\mathcal{V}(\phi)=\text{const.}} \Lambda_{F,s}[\phi],$$

where $\mathcal{V}(\phi)$ denotes the Lebesgue measure of $\phi(\Omega)$, i.e.,

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$$\operatorname{Ker} d\mathcal{V}(\phi) \subseteq \operatorname{Ker} d\Lambda_{F,s}[\phi].$$



Let
$$\mathcal{V}_0 > 0$$
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Theorem

Let Ω be a bounded domain of \mathbb{R}^N of class C^1 . Let $\tilde{\phi}$ be such that $\tilde{\phi}(\Omega)$ is a ball. Let $\tilde{\lambda}$ be an eigenvalue of problem (1) in $\tilde{\phi}(\Omega)$, and let F be the set of $j \in \mathbb{N}$ such that $\lambda_j[\tilde{\phi}] = \tilde{\lambda}$. Then $\Lambda_{F,s}$ has a critical point at $\tilde{\phi}$ on $V(\mathcal{V}(\tilde{\phi}))$, for all s = 1, ..., |F|.





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Hence, balls are critical domains for all simple eigenvalues and for all the symmetric functions of all multiple eigenvalues under measure constraint.



Can we say more on the critical nature of balls for the Steklov eigenvalues?



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Theorem

Among all bounded domains of class C^1 with fixed measure, the ball is the unique maximizer of the first non-negative eigenvalue of problem (1), that is

 $\lambda_2(\Omega) \leq \lambda_2(\Omega^*),$

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- Use of suitable test functions built from the eigenfunctions of the ball in the variational characterization of the eigenvalues.
- (Classical) isoperimetric inequality for weighted perimeters.



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- "If Ω is such that λ₂(Ω) ~ λ₂(Ω*), then Ω has to resemble a ball?"
- "In which way this is quantified?"



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$$\mathcal{A}(\Omega) = \inf \left\{ \frac{|\Omega \triangle B|}{|\Omega|} : B \text{ ball with } |B| = |\Omega| \right\}$$

be the so-called Fraenkel Asymmetry.





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be the so-called Fraenkel Asymmetry.

Fraenkel Asymmetry measures the "distance" in the L^1 sense of a generic set from the "family" of balls.





We have the following

Theorem

For every domain Ω in \mathbb{R}^N of class C^1 the following estimate holds:

$$\lambda_{2}(\Omega) \leq \lambda_{2}(\Omega^{*}) \left(1 - c_{N} \mathcal{A}(\Omega)^{2}\right),$$
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where c_N is a suitable constant and Ω^* is a ball with the same measure as Ω .



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This is the isoperimetric inequality in quantitative form.



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To do so we shall exhibit a family $\{\Omega_{\varepsilon}\}$ of sets approaching the unit ball B such that

$$\mathcal{A}(\Omega_{\varepsilon}) \simeq \frac{|\Omega_{\varepsilon} \triangle B|}{|\Omega_{\varepsilon}|} \simeq \varepsilon \text{ and } \lambda_2(B) - \lambda_2(\Omega_{\varepsilon}) \simeq \varepsilon^2, \quad \varepsilon \ll 1.$$









Classical isoperimetric inequality: if E is a Borel set in R^N with finite Lebesgue measure, then the ball with the same measure has lower perimeter, that is

 $P(E) \ge P(\Omega^*),$

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N. Fusco, F. Maggi, and A. Pratelli. The sharp quantitative isoperimetric inequality. *Ann. of Math. (2) 168* (2008), no. 3, 941–980.



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L. Brasco, G. De Philippis, and B. Velichkov. Faber-Krahn inequalities in sharp quantitative form. *Duke Math. J. 164* (2015), no. 9, 1777–1831.



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This means that $\mathcal{A}(\Omega_{\varepsilon}) \simeq \frac{|\Omega_{\varepsilon} \triangle B|}{|\Omega_{\varepsilon}|} \simeq \varepsilon$ and moreover

 $P(\Omega_{\varepsilon}) - P(B) \simeq \varepsilon^2$ and $\mu_1(\Omega_{\varepsilon}) - \mu_1(B) \simeq \varepsilon^2 \quad \varepsilon \ll 1.$



In the case of the first positive eigenvalue of our biharmonic Steklov problem, we have that

$$\lambda_2(B) - \lambda_2(\Omega_{\varepsilon}) \simeq \varepsilon, \quad \varepsilon \ll 1,$$

when Ω_{ε} are nearly spherical ellipsoids.



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Why? Is 1 the right exponent for the Fraenkel asymmetry in the isoperimetric inequality or there exist suitable families $\{\Omega_{\varepsilon}\}$ such that the rate of convergence is ε^2 , proving the sharpness of the exponent 2?

The answer relies not only on the geometry of the critical set (the ball), but also on the features of the problem itself (that is, on the properties of the eigenfunctions).



Some observations:

An eigenvalue does not have a straightforward geometrical meaning like in the case of the perimeter, for example. Thus, it is not straightforward to understand how deformations of an optimal shape affect the eigenvalues.



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- If an eigenvalue is "shape differentiable" (e.g., if it is always simple as in the case of μ₁(Ω)), then its "shape derivative" at the "extremal point" would be zero (this is rather heuristic).



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- An eigenvalue does not have a straightforward geometrical meaning like in the case of the perimeter, for example. Thus, it is not straightforward to understand how deformations of an optimal shape affect the eigenvalues.
- If an eigenvalue is "shape differentiable" (e.g., if it is always simple as in the case of μ₁(Ω)), then its "shape derivative" at the "extremal point" would be zero (this is rather heuristic).
- In the case of μ₁(Ω) (first Dirichlet eigenvalue), any perturbation Ω_ε = (Id + εV)B, for some smooth vector field V, should provide an expansion of the form

$$\mu_1(\Omega_{\varepsilon}) \simeq \mu_1(B) + O(\varepsilon^2), \quad \varepsilon \ll 1.$$



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To prove that the exponent is sharp, we have to **exclude** that this happens for every "direction".



We define a family $\{\Omega_{\varepsilon}\}$ in this way

$$\Omega_{\varepsilon} = \left\{ x \in \mathbb{R}^N : |x| < 1 + \varepsilon \psi(x/|x|) \right\},\$$

where $\psi \in C^{\infty}(\partial B)$





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This family of sets is such that $\mathcal{A}(\Omega_{\varepsilon}) \simeq \varepsilon$ and $\lambda_2(B) - \lambda_2(\Omega_{\varepsilon}) \simeq \varepsilon^2$, proving that the exponent 2 is sharp.









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- **3** $\int_{\partial B} (a \cdot x)^2 \psi d\sigma = 0$ for all $a \in \mathbb{R}^N$ has a strict relation with the problem: any eigenfunction ξ corresponding to $\lambda_2(B)$ is of the form $(a \cdot x)$ for some $a \in \mathbb{R}^N$



- 1 $\int_{\partial B} \psi d\sigma = 0$ has a pure geometrical meaning: this implies that Ω_{ε} has the same measure as *B* up to an error of ε^2 ;
- 2 ∫_{∂B}(a · x)ψdσ = 0 for all a ∈ ℝ^N has again a pure geometrical meaning: this implies that the barycenter of Ω_ε is the origin up to an error of ε². In particular this also implies that A(Ω_ε) ≃ ε,

3 $\int_{\partial B} (a \cdot x)^2 \psi d\sigma = 0$ for all $a \in \mathbb{R}^N$ has a strict relation with the problem: any eigenfunction ξ corresponding to $\lambda_2(B)$ is of the form $(a \cdot x)$ for some $a \in \mathbb{R}^N$, so that

$$\left|\xi\right|^2_{\mid_{\partial B}} = (a\cdot x)^2 \quad \text{and} \quad \left|D^2\xi\right|^2 + \tau \left|D\xi\right|^2_{\mid_{\partial B}} = (b\cdot x)^2,$$

for some $a, b \in \mathbb{R}^N$ and such relations are crucial in proving $\lambda_2(B) - \lambda_2(\Omega_{\varepsilon}) \simeq \varepsilon^2$ for $\varepsilon \ll 1$.



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Along the "ellipsoid" direction, $\lambda_2(\Omega)$ has a non-trivial super-differential



First two eigenvalues of nearly spherical ellipsoids Ω_{ε} and $\tau = 1$.





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First two (multiple) eigenvalues of the "flower domains" Ω_{ε} and $\tau=1.$





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First two eigenvalues of Ω_{ε}



Note that also in this case the eigenvalue is not differentiable at the "maximum point" but converges with the sharp exponent 2.



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THANK YOU